

# Preliminary studies of Energetic Particle driven Alfvénic modes for a model DTT equilibrium using HYMAGYC

G. Vlad, S. Briguglio, G. Fogaccia and V. Fusco  
ENEA – CR Frascati (Rome) ITALY

**CNPS Workshop on Multi-scale Energetic particle Transport in fusion  
devices Jointly organized with the  
MET Project Final Workshop  
March 3 - 5, 2021**

- Introduction
- HYMAGYC & case study for DTT
- MHD characterization of the equilibrium
- Energetic particle driven Alfvénic modes vs. toroidal mode number
- Conclusions

DTT, the Divertor Tokamak Test facility, is the new plasma physics research device under construction in Italy, which will benefit from a substantial support from EUROfusion to specifically address the problem of heating and power exhaust in ITER and DEMO devices. DTT characteristic parameters considered here are: toroidal field  $B_0 \approx 6.0\text{T}$ , major radius  $R_0 \approx 2.08\text{m}$ , aspect ratio  $A \approx 3.2$ , plasma current  $I_p \approx 5.5 \text{ MA}$ , additional power  $P_{\text{Tot}} \approx 45 \text{ MW}$  (**note that the device parameters have been slightly updated, e.g.,  $R_0 \approx 2.18\text{m}$** ).

The equilibrium and plasma parameters for the Single Null (SN) original baseline scenario (2018) can be found at [https://www.afs.enea.it/vlad/Miscellaneous/Testcase\\_DTT\\_HYMAGYC\\_2019/](https://www.afs.enea.it/vlad/Miscellaneous/Testcase_DTT_HYMAGYC_2019/):

It is a “engineering” equilibrium, with bulk ion density resulting from a METIS simulation: energetic particle density profile as from ITER-SC2 scenario; high resolution model equilibrium recomputed by CHEASE ( $B_0, I_p > 0$ , COCOS=2 coordinate system)

DTT, v1, Single Null, t=36s

$B_0 \approx 6.0\text{T}$

$R_0 \approx 2.08\text{m}$

$\beta_{\text{th},0} \approx 4\%$

$T_H = 0.45\text{MeV}$

$n_{H0}/n_{i0} \approx 0.05$

$\beta_{H0}/\beta_{\text{th},0} \gtrsim 1$

$m_H/m_i = 0.5$  (H EPs/D bulk ions)

$\rho_H/a \approx 0.017$

$v_H/v_{A0} \approx 1$ .

$\Gamma = 5/3$  (adiabatic coefficient)

$\eta = 0$  (ideal plasma)

Isotropic Maxwellian EPs, FOW only

Quantity	Value	Data definition/Origin
$B_{\text{geo}}$ [T]	5.9994	EQDSK, vacuum magnetic field at $R=R_{\text{geo}}$
$R_{\text{geo}}$ [m]	2.0802	geometric major radius $(R_{\text{LCMS\_max}}+R_{\text{LCMS\_min}})/2$
$B_0$ [T]	$B_0=B_{\text{geo}}$	normalization coefficient for the magnetic field
$R_0$ [m]	$R_0=R_{\text{geo}}$	normalization coefficient for the lengths
$a$ [m]	0.65	minor radius $(R_{\text{LCMS\_max}}-R_{\text{LCMS\_min}})/2$
$\epsilon_{\text{dev}}$ [m]	0.31247	inverse Aspect ratio $(a/R_{\text{geo}})$
$n_{i0}$ [ $10^{20}/\text{m}^3$ ]	2.0739	from METIS simulation
$n_{\text{EP0}}/n_{i0}$	0.05	EP density/bulk ion density
$m_i/Z_i$	2/1	bulk ion mass/charge (D) (in units of proton mass/electron charge)
$m_{\text{EP}}/Z_{\text{EP}}$	1/1	EP mass/charge (H) (in units of proton mass/electron charge)
$m_{\text{EP}}/m_i$	0.5	mass ratio (EP/bulk ion)
$T_{\text{EP0}}$ [MeV]	0.45	on-axis EP Temperature (constant on radius), Maxwellian distribution
$v_{A0}$ [m/s]	$6.42242 \times 10^6$	on-axis Alfvén velocity $\Rightarrow 2.18 \times 10^6 B_0[\text{T}] / \sqrt{m_i n_{i0} [10^{20}/\text{m}^3]}$
$\tau_{A0}$ [s]	$3.23865 \times 10^{-7}$	$R_0/v_{A0}$
$\omega_{A0}$ [rad/s]	$3.08770 \times 10^6$	$1/\tau_{A0}$
$v_{\text{EPth0}}$ [m/s]	$6.92258 \times 10^6$	$\sqrt{T_{\text{EP0}}/m_{\text{EP}}} \Rightarrow 9.79 \times 10^6 \sqrt{T_{\text{EP0}}[\text{MeV}]/m_{\text{EP}}}$ note the definition w/o $\sqrt{2}$ !
$v_{\text{EPth0}}/v_{A0}$	1.07788	
$\omega_{ci}$ [rad/s]	$5.747425 \times 10^8$	EP gyrofrequency $\Rightarrow 9.58 \times 10^7 Z_{\text{EP}} B_0[\text{T}] / m_{\text{EP}}$
$\rho_{\text{EP0}}$ [m]	0.01204466	on-axis EP Larmor radius $(v_{\text{EPth0}}/\omega_{ci}) \Rightarrow 0.102 \sqrt{m_{\text{EP}} T_{\text{EP0}}[\text{MeV}]} / (Z_{\text{EP}} B_{\text{mag}}[\text{T}])$
$\rho_{\text{EP0}}/R_0$	0.00577926	on-axis EP Larmor radius/ $R_0$
$\rho_{\text{EP0}}/a$	0.0184954	on-axis EP Larmor radius/ $a$

# HYMAGYC code description

HYMAGYC [1] is a hybrid MHD-Gyrokinetic code.

- It is suited to study the interaction between EPs and Alfvénic modes, for high- $\beta$  axisymmetric equilibria
- Electromagnetic fields are fully retained: electrostatic potential  $\varphi$  and vector potential  $\mathbf{A}$
- Thermal plasma is described as a single fluid by **full**, resistive, linear MHD equations.
- The fields solver originates from MARS [2], transformed from an eigenvalue solver to an initial value one.
- Energetic particles are described by nonlinear gyrokinetic Vlasov equations [3] expanded up to order  $O(\varepsilon^2)$  and  $O(\varepsilon\varepsilon_B)$  and solved by particle-in-cell (PIC) techniques.
- The MHD and the gyrokinetic modules, are coupled together by inserting the divergence of the EP pressure tensor in the MHD momentum equations [4]
- It uses the equilibrium solver CHEASE; it is IMAS compliant

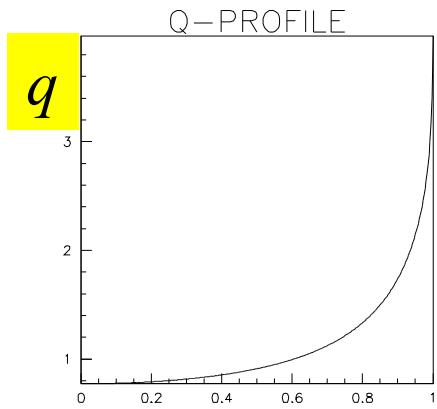
## ORDERING and DEFINITIONS

- gyrokinetic ordering parameter  $\varepsilon \approx \rho_H/L_n$
- $\varepsilon_B \approx \rho_H/L_B$ ,
- $\rho_H$  the EP Larmor radius
- $L_n/L_B$  the characteristic length scales of the equilibrium plasma density/magnetic field.
- Space-time ordering for the fluctuating electromagnetic fields:  $k_\perp \rho_H = O(1)$ ,  $k_\parallel \rho_H = O(\varepsilon)$ ,  $\omega/\Omega_H = O(\varepsilon)$
- $k_\perp$  the perpendicular (to the equilibrium magnetic field) wave vector
- $k_\parallel$  the parallel one
- $\omega$ : characteristic fluctuation frequency and  $\Omega_H$  the EP gyrofrequency.

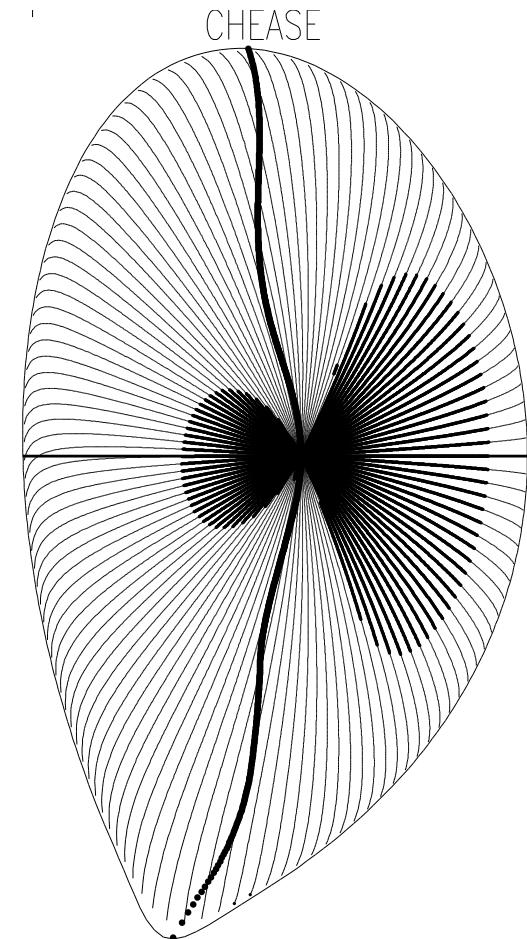
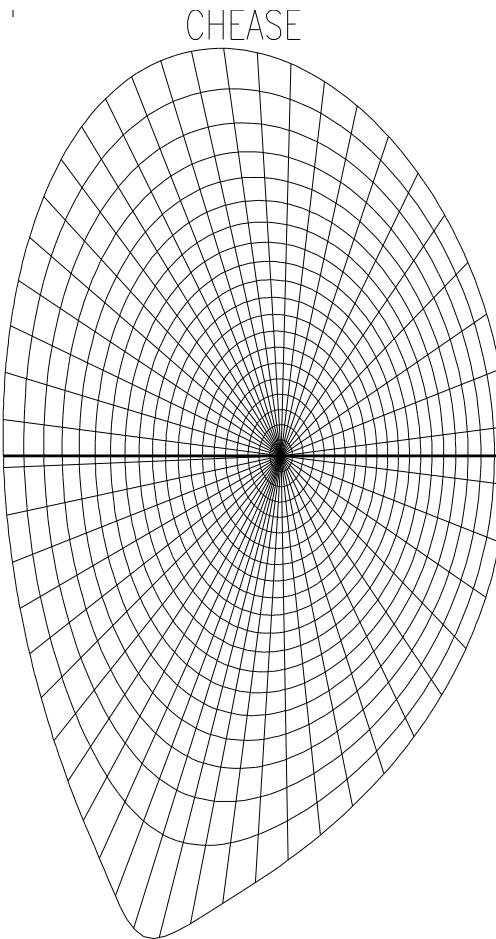
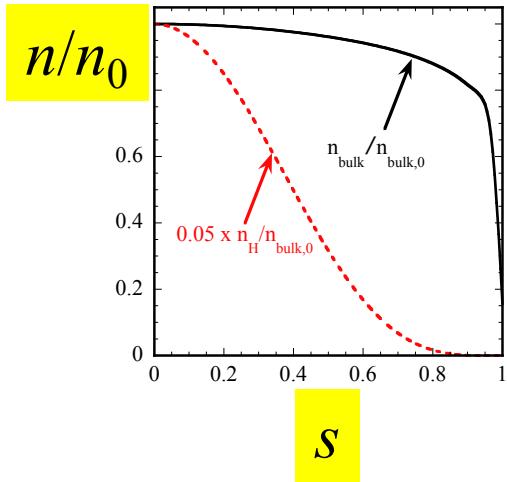
## References

- [1] G. Fogaccia, G. Vlad, S. Briguglio, Nucl. Fusion 56 (2016) 112004
- [2] Bondeson A., Vlad G. and Lütjens H. 1992 IAEA Technical Committee Meeting on Advances in Simulations and Modelling of Thermonuclear Plasmas (Montreal, 15–17 June 1992) p. 306 (Vienna, Austria: International Atomic Energy Agency)
- [3] Brizard A.J. and Hahm T.S. 2007 Rev. Mod. Phys. 79 421–68
- [4] Park W. et al 1992 Phys. Fluids B 4 2033
- [5] T. Wang, Z. Qiu, F. Zonca, S. Briguglio, G. Fogaccia, G. Vlad, and X. Wang, Phys. Plasmas 25, 062509 (2018)
- [6] T. Wang, X. Wang, S. Briguglio, Z. Qiu, G. Vlad, and F. Zonca, Physics of Plasmas 26, 012504 (2019)

# DTT equilibrium (CHEASE), reference case



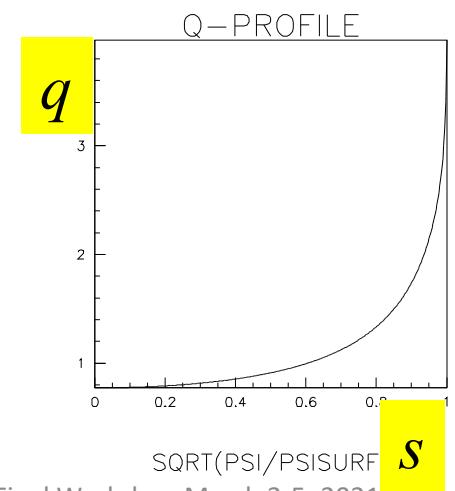
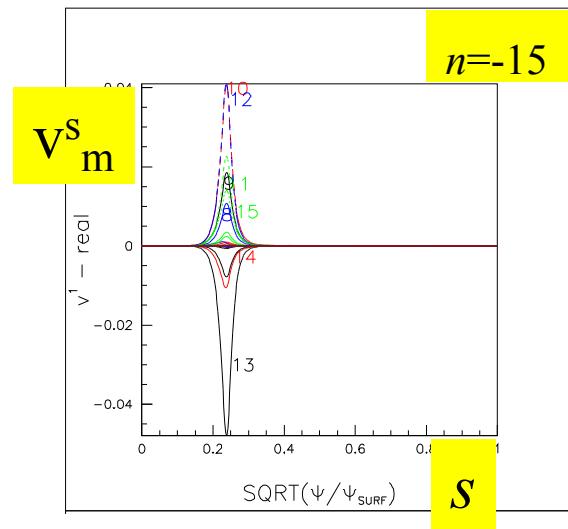
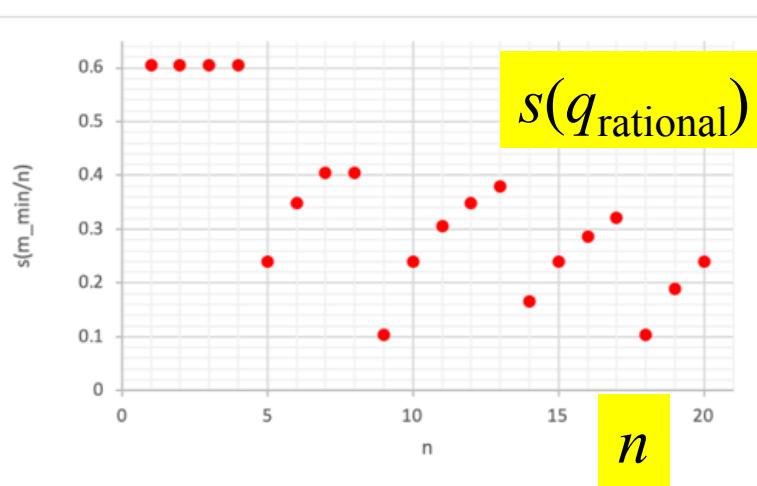
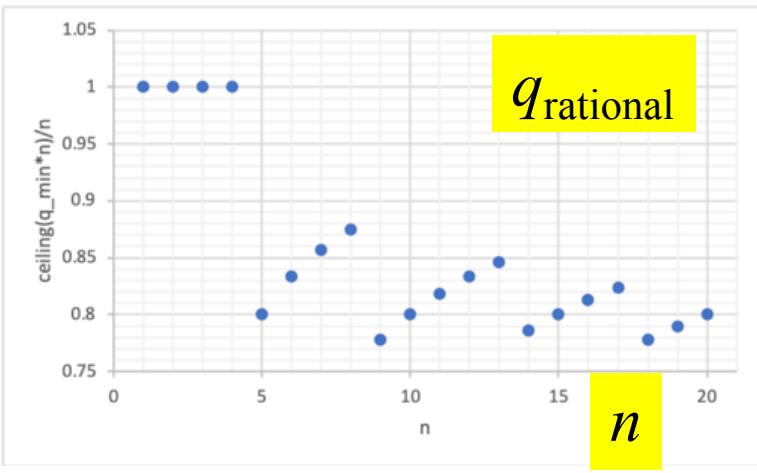
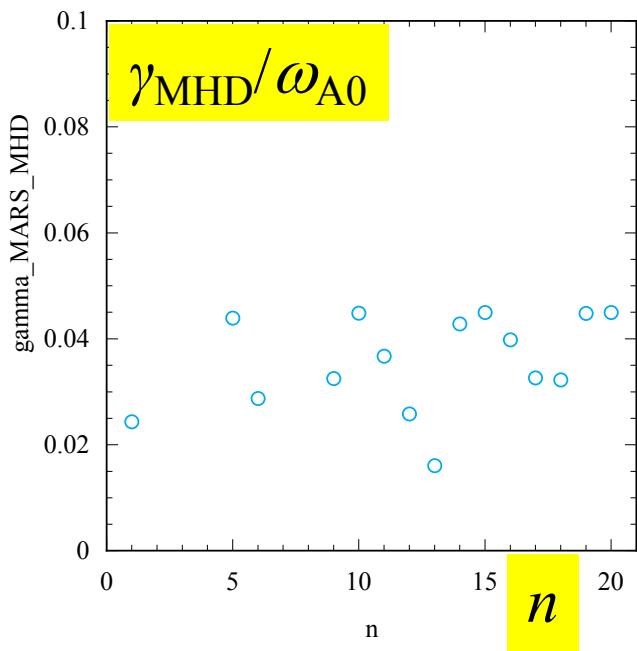
$$s \equiv \sqrt{\psi/\psi_{\text{surf}}}$$



```
***  
*** EquilzDTTzSNzv1z|plz5d5MAzR0z2d08zbetapz0d43zliz0d895ztatz36szrefzmesh-rwle  
*** cocoszin=11, ntmf0=0, ncscal=4 (no scaling), tensbnd=-300.0,  
*** cocoszout=13,
```

# MHD unstable modes, $\omega/\omega_{A0}=0$ : infernal modes (MARS)

*"When the shear is sufficiently weak, the oscillations can result in bands of unstable n-values, which are present even when the standard ballooning theory predicts complete stability. These instabilities are named 'infernal modes'. The occurrence of these instabilities at integer n is shown to be a sensitive function of the q-axis, raising the possibility of a sharp onset as the plasma parameters evolve."* (from the original paper by J. Manickam et al 1987 Nucl. Fusion 27 1461)



# Adding Energetic Particle contribution

- Linear growth-rate of Alfvén modes and EPMS is proportional to the energetic particle diamagnetic drift frequency  $\omega_{*Es}$  (Liu Chen et al, PRL 1984);
- from gyrokinetic description of the perturbed EP distribution function, stability given by  $\delta W_k$  trapped particles:

$$\widehat{\delta W}_k \propto \mathcal{I} \circ \frac{1}{n\bar{\omega}_d + \ell\omega_b - \omega} QF_0$$

circulating particles:  $\widehat{\delta W}_k \propto \mathcal{I} \circ \frac{1}{n\bar{\omega}_d + [\ell + nq(r) - m]\omega_t - \omega} QF_0$

$\mathcal{I} \circ$  : integro-differential operator

$$QF_0 = \omega \left( \partial_E + \frac{\widehat{\omega}_{*Es}}{\omega} \right) F_0$$

$$\widehat{\omega}_{*Es} F_0 = \frac{1}{\omega_c} \frac{\mathbf{k} \times \mathbf{B}}{B} \cdot \nabla F_0$$

$$\omega_{*Es} \approx \frac{m_{pol}}{n_{He} e_s r B} \frac{dp_H}{dr} = \frac{nq(r)}{n_{He} e_s r B} \frac{dp_H}{dr}$$

$\omega_{*Es} \propto nq(r)$ : short wavelength favored

- Lower bound for wavelength set by characteristic energetic particle orbit width,  $\rho_E$ :  $k_\perp \rho_E \approx 1$  (finite orbit width averaging; magnetic drifts usually larger than Larmor radius)
- optimal condition:  $n_{max} q \lesssim (r/\rho_E)$
- $n_{max-ITER} \approx O(10)$

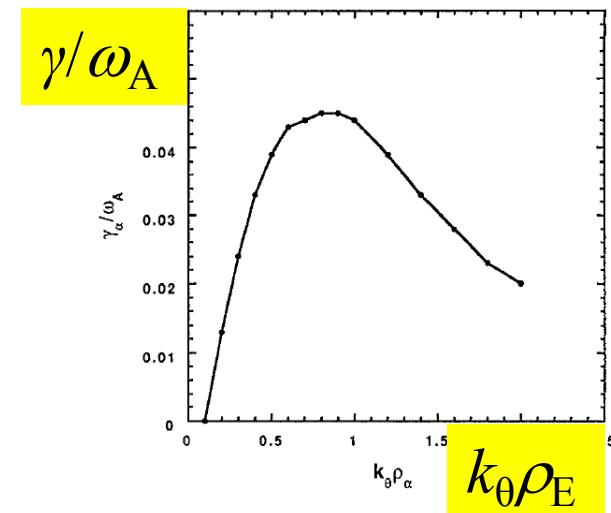


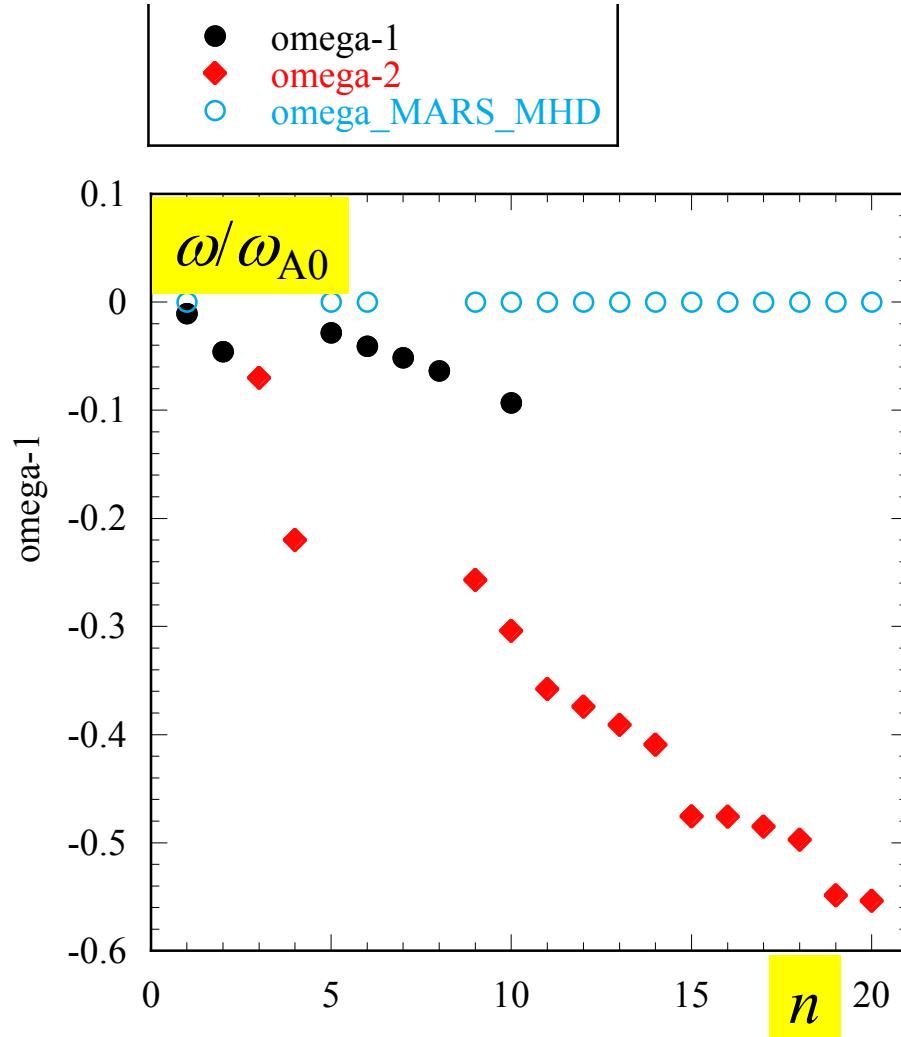
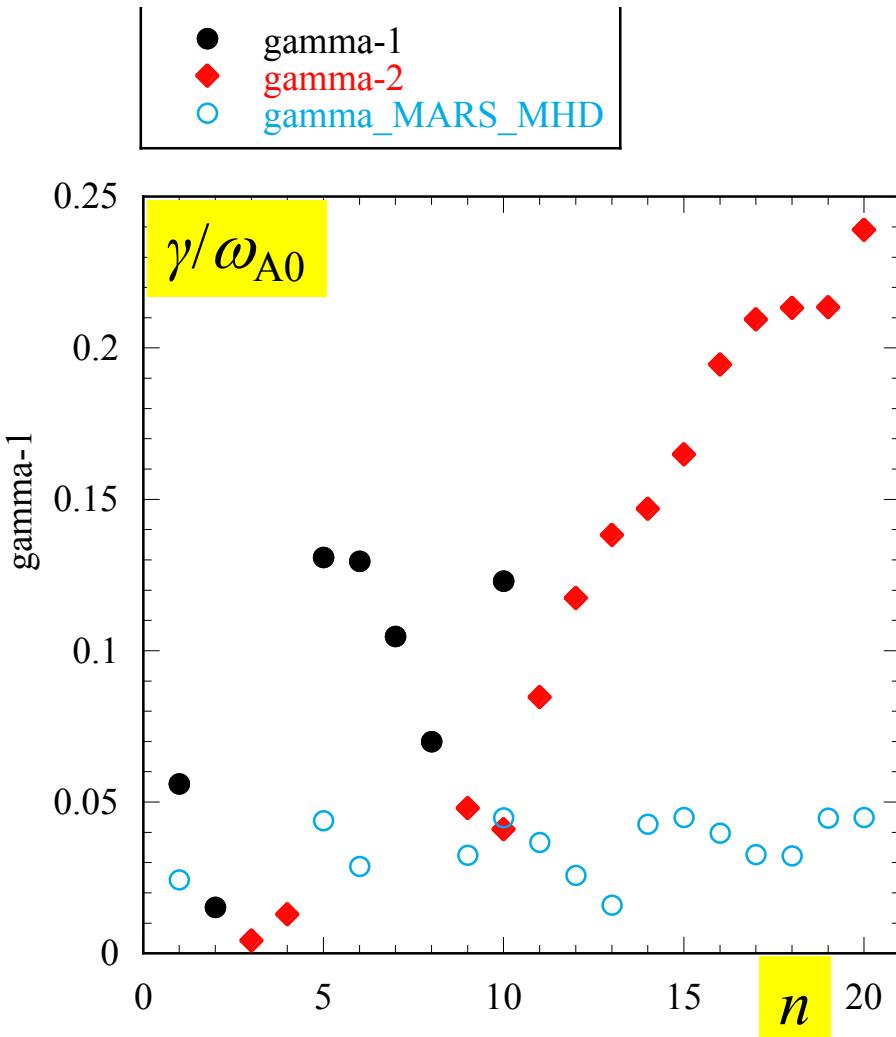
FIG. 3. Growth rate induced by circulating alpha particles as a function of  $k_\theta \rho_\alpha$  for  $s=0.6$ ,  $\Delta_p=0$ , and  $v_a/v_A=2.0$ .

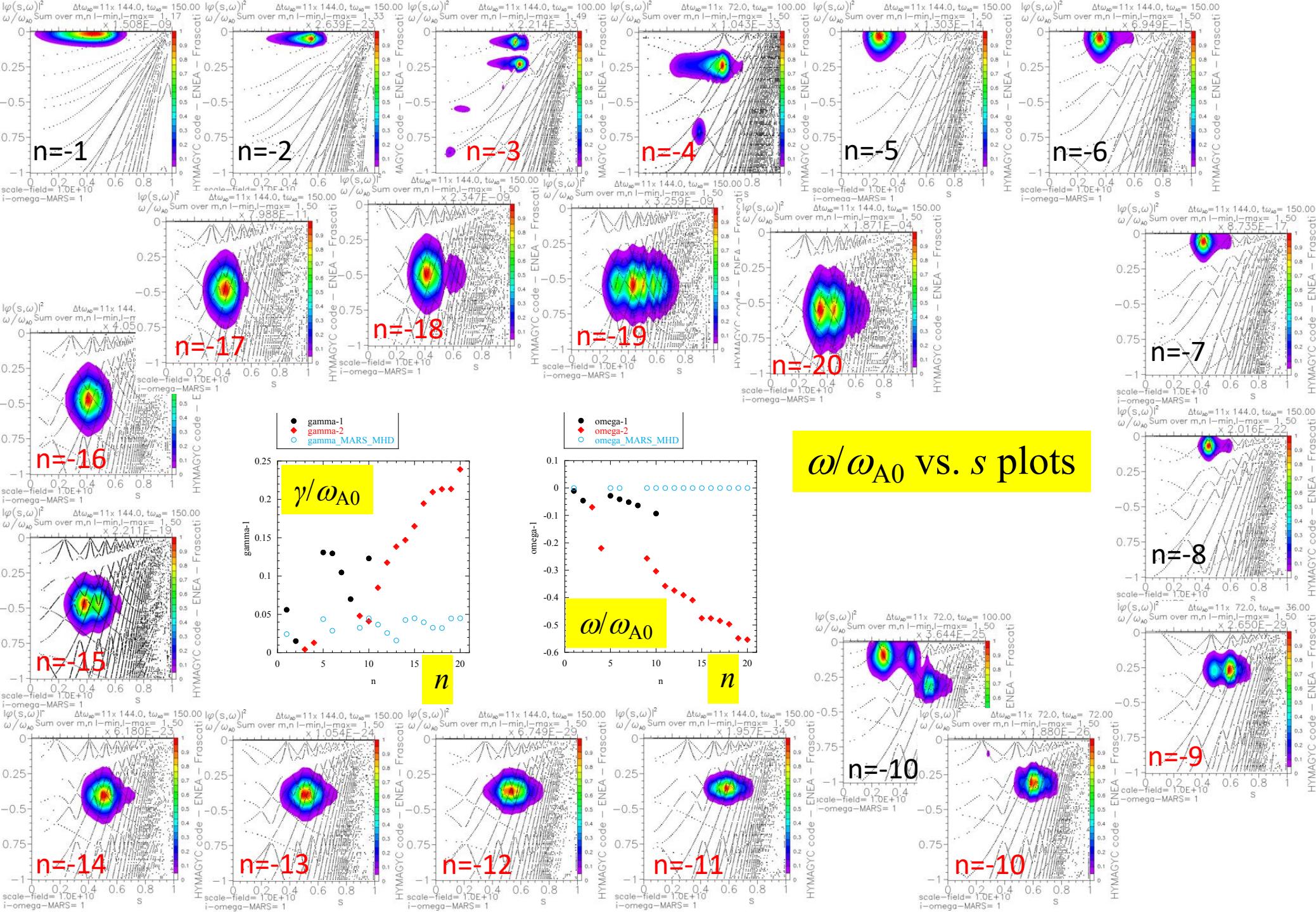
## Adding Energetic Particles contribution (contd.)

- $\gamma_{\text{MHD}}, \omega_{\text{MHD}}$ : "infernal" modes

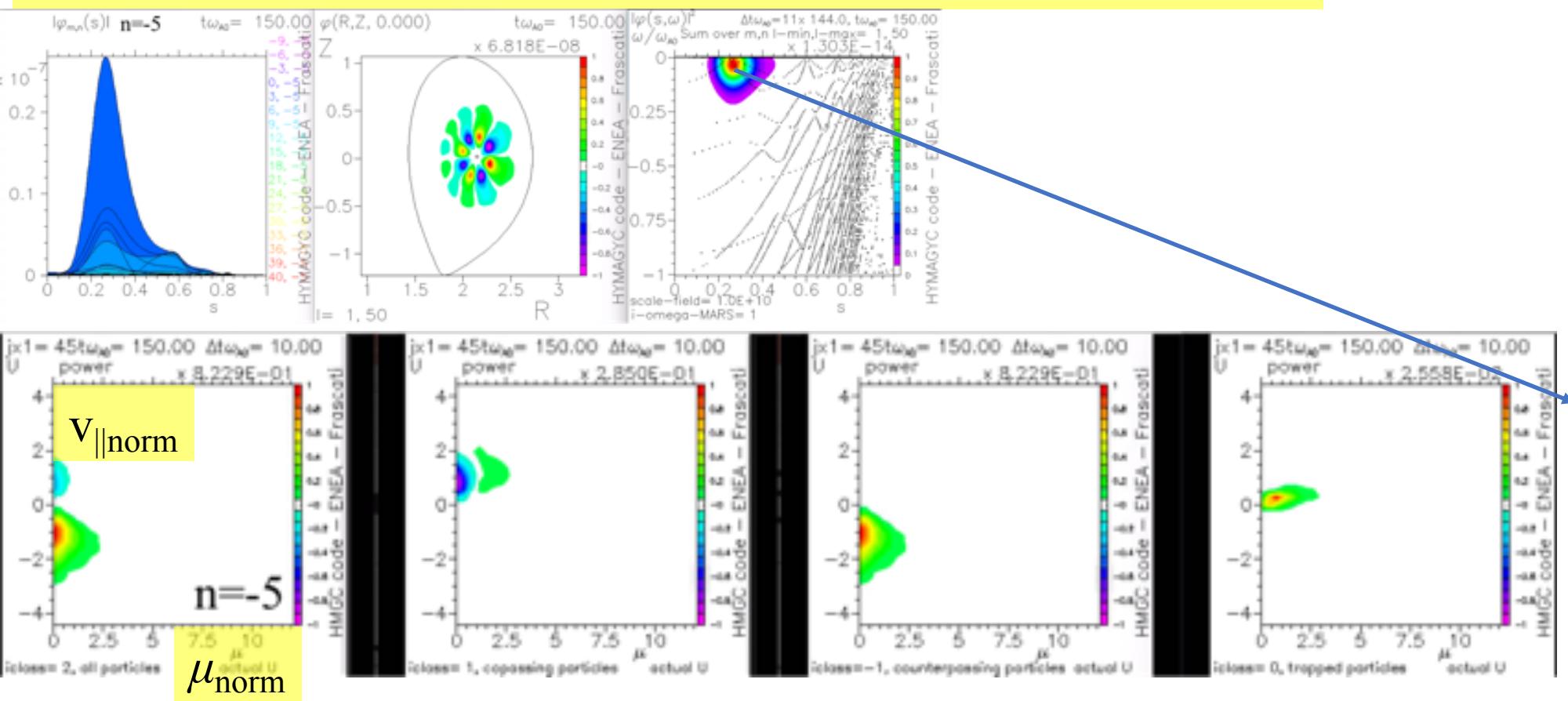
Adding Energetic Particle contribution:

- $\gamma_1, \omega_1$ : EP driven "infernal" modes
- $\gamma_2, \omega_2$ : EP driven Alfvénic modes

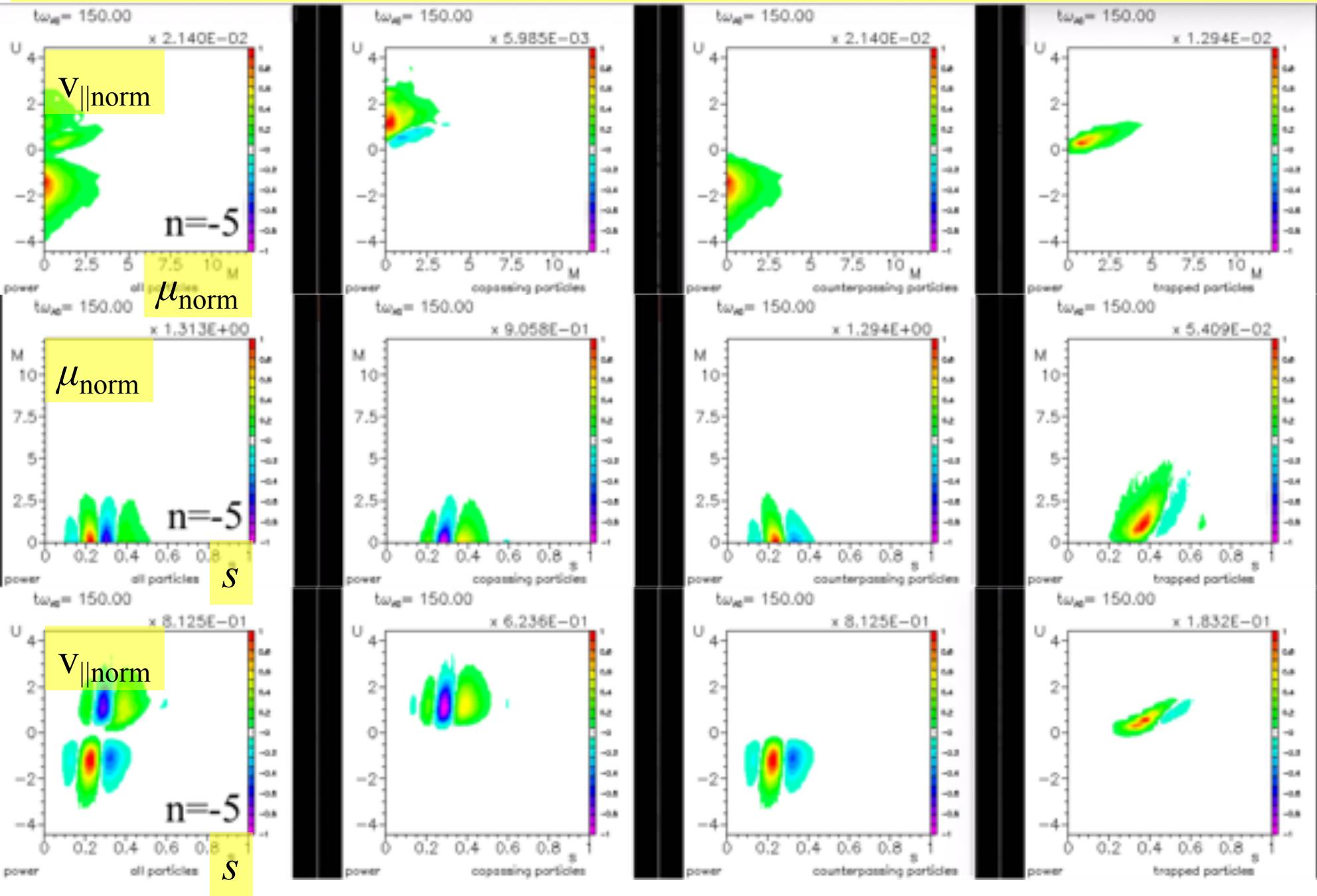




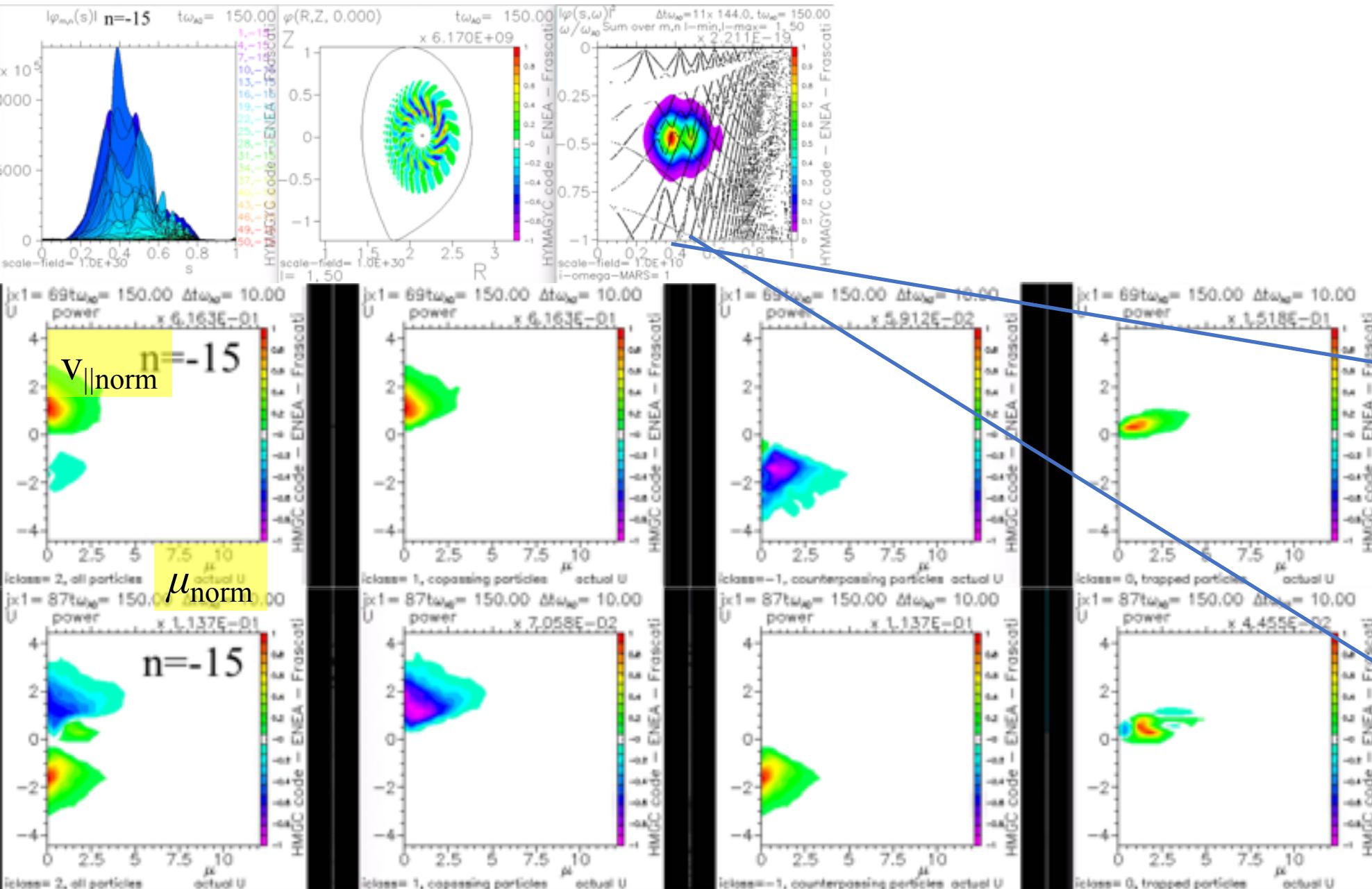
# Power transfer between EPs and wave P(s,M,U); n=-5 – 3d plots => counter-passing



Power transfer between EPs and wave P(s,M,U); n=-5 – 2d plots (reduced over third coord.) => counter-passing



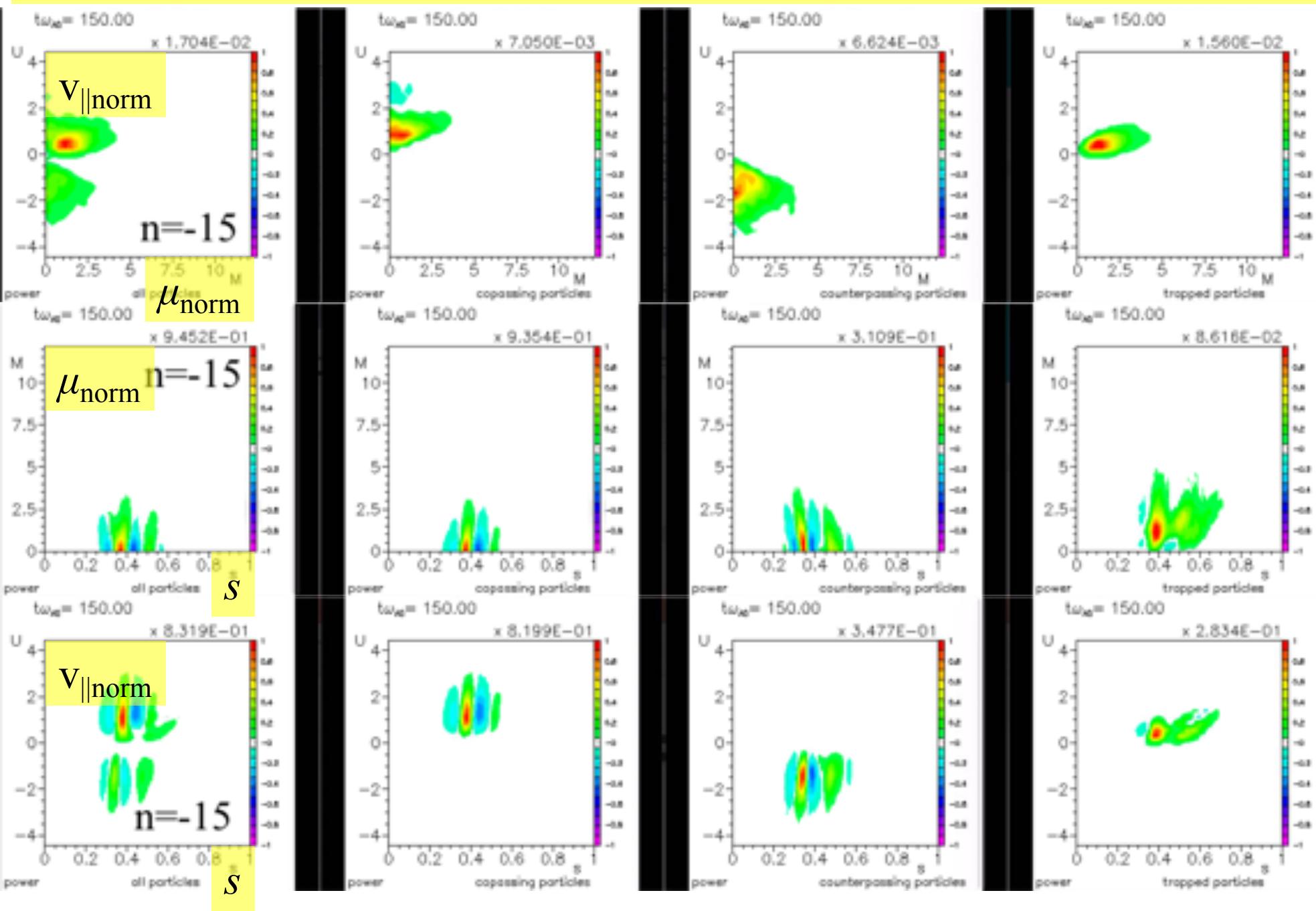
# Power transfer between EPs and wave P(s,M,U); n=-15 – 3d plots => trapped particles



$s \approx 0.38$  ( $jx1=69$ )

$s \approx 0.48$  ( $jx1=87$ )

Power transfer between EPs and wave  $P(s, M, U)$ ;  $n = -15$  – 2d plots (reduced over third coord.) => trapped particles

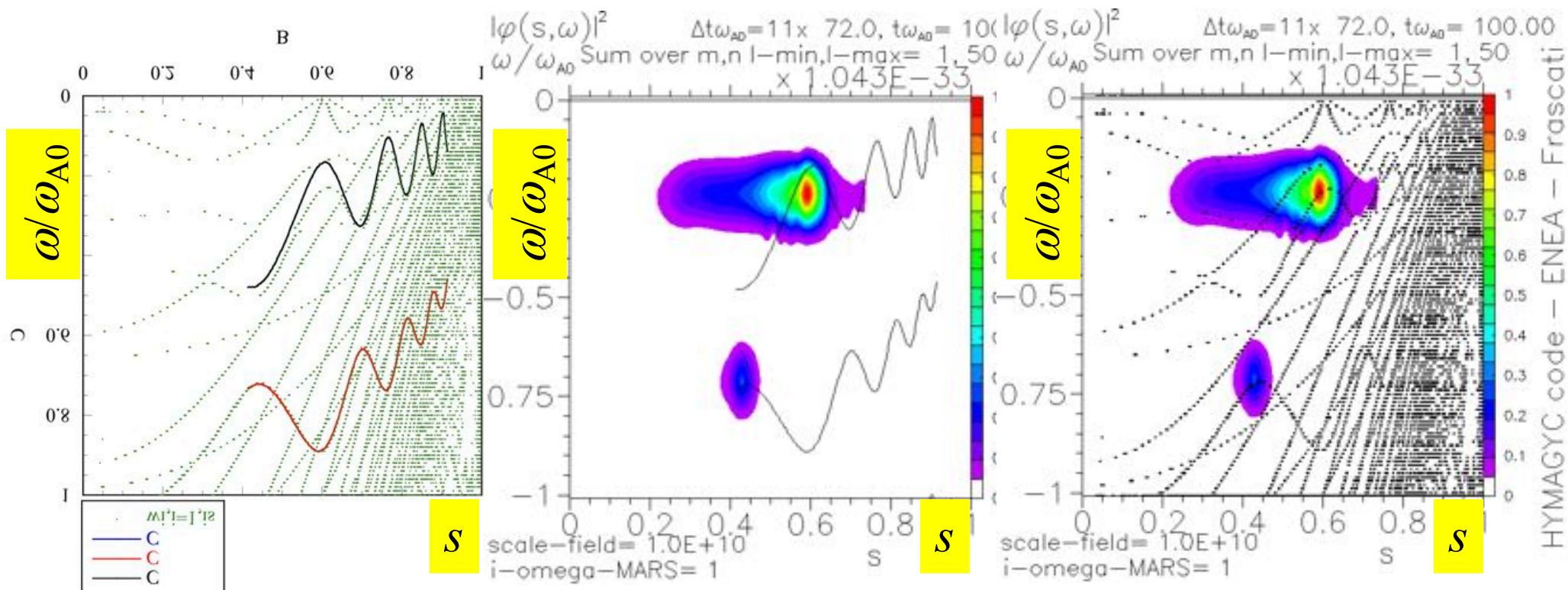


## $n=-4$ Alfvén continuous spectrum, detailed structure

comparison between MARS continua and approximated (Falcon) Shear Alfvén Wave continua using slow-sound approximation: comparison is satisfactory for the upper continuum of the toroidal gap, not satisfactory for the lower continuum and below (BAAE gap, etc.).

GREEN: complete MARS continua;

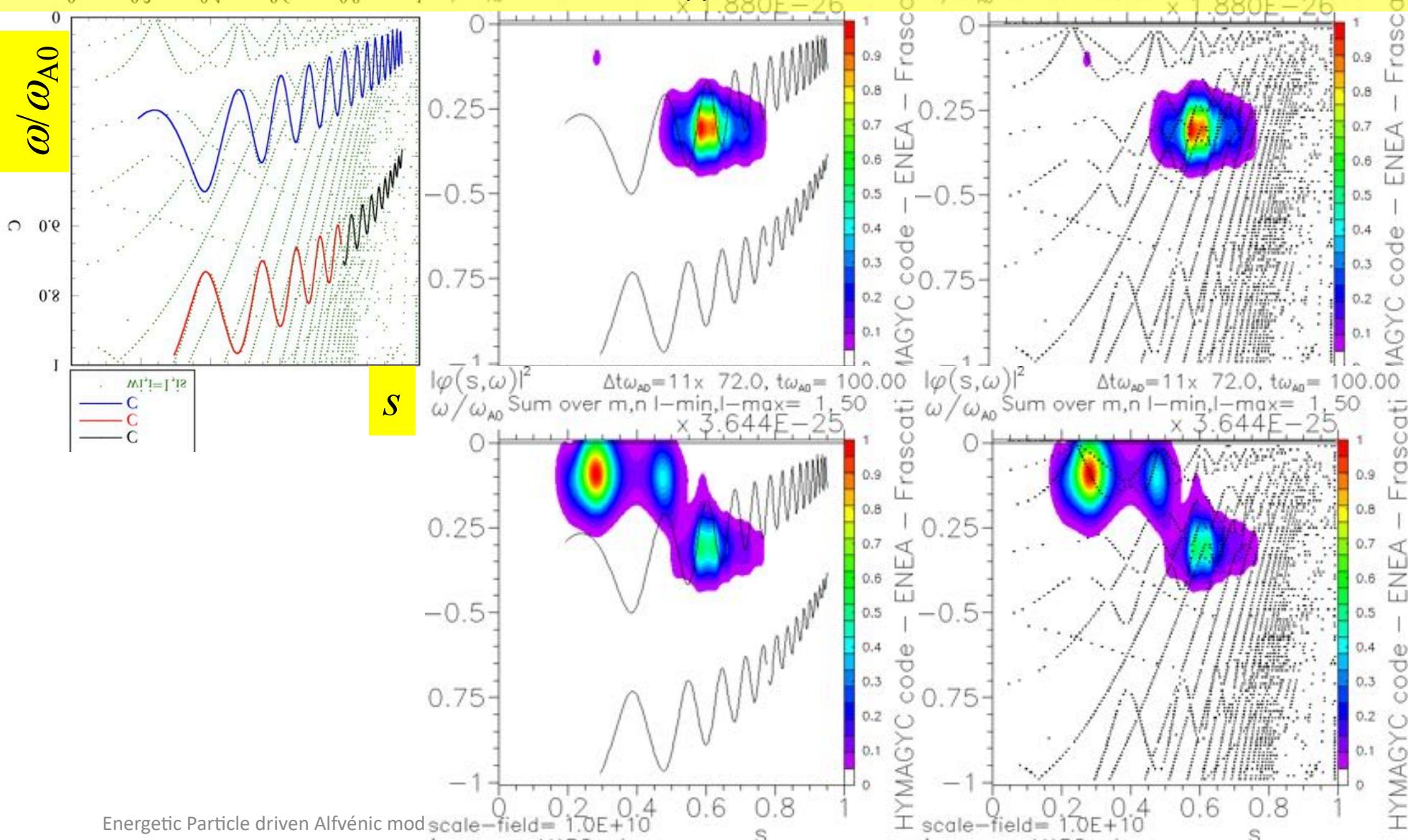
BLACK, RED solid lines: Falcon with slow sound approximation



**n=-10 Alfvén continuous spectrum**, detailed structure: comparison between MARS continua and approximated (Falcon) Shear Alfvén Wave continua using slow-sound approximation: comparison is satisfactory for the upper continuum of the toroidal gap, not satisfactory for the lower continuum and below (BAAE gap, etc.).

GREEN: complete MARS continua;

BLACK, BLUE, RED solid lines: Falcon with slow sound approximation

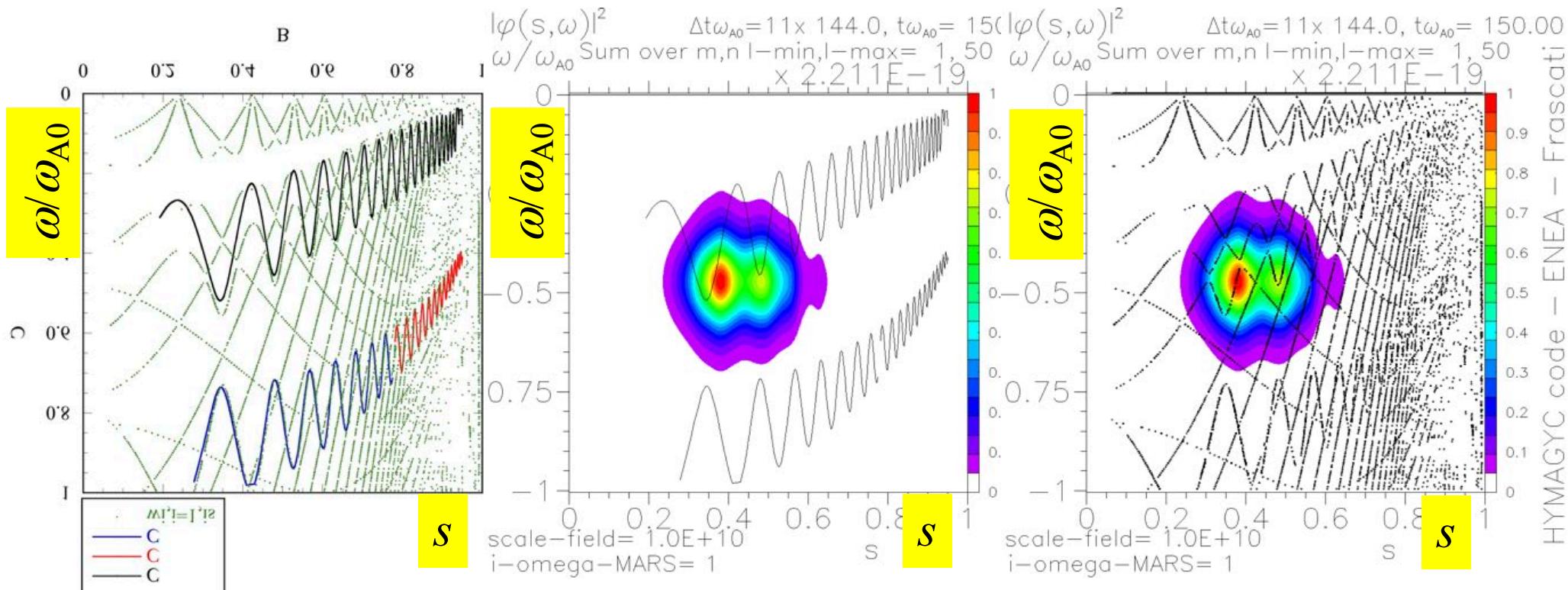


## $n=-15$ Alfvén continuous spectrum, detailed structure

comparison between MARS continua and approximated (Falcon) Shear Alfvén Wave continua using slow-sound approximation: comparison is satisfactory for the upper continuum of the toroidal gap, not satisfactory for the lower continuum and below (BAAE gap, etc.).

GREEN: complete MARS continua;

BLACK, BLUE, RED solid lines: Falcon with slow sound approximation



# Conclusions

- First systematic investigation on Energetic Particle driven Alfvénic Modes on DTT model equilibrium using HYMAGYC:
  - useful exercise to test HYMAGYC potentialities in realistic scenarios environment;
  - moderately high toroidal mode numbers have been tested ( $n \leq 20$ );
  - useful testbed for checking numerics, convergence, etc.
- Model equilibrium should be updated to latest DTT scenarios (slightly larger device dimensions, different heating mix, use proper  $B_0, I_p$  signs, etc.), possibly using more realistic scenarios as obtained by e.g., transport code simulations (consistent scenario...)
- More appropriate Energetic Particles distribution function to be used (e.g., anisotropic slowing down, instead of isotropic Maxwellian, Constant of Motion parametrized distribution function, etc.)
- Hamiltonian mapping “machinery”, largely used with HMGC, is available also for HYMAGYC
- ...

