

# Alfvén wave simulation using implicit particle scheme: from 1D to torus geometry

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# Ongoing work for MET project

1. Development of full $f$ , implicit, mixed Particle-in-Cell-Particle-in-Fourier (PIC-PIF) method for simulation of Alfvén wave and EP physics (2019,2020)	WP2-WP3
2. Anisotropic EP effects on zonal flow residual (mainly in 2019; paper to be written in 2020)	WP1

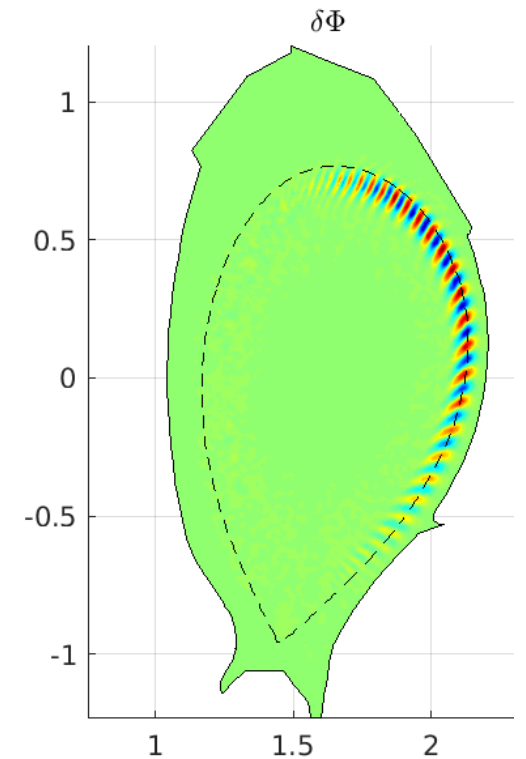
This talk is for Part. 1

# Background and motivation

- Related work on implicit scheme

[Günter JCP09, Kleiber POP16, Perse, Numkin19, ELMFIRE, XGC]

- Implicit scheme can be complementary to CT, pullback...
  - For a system with  $N_G$  field grids and  $N_p$  particles, the total Degree of Freedom (DOF) is  $(N_G + N_p)$
  - The direct implicit solver requires solving a linear system with  $\text{DOF} = N_G + N_p$ . Matrix size is  $(N_G + N_p)^2$  !!!
  - Solution: “particle enslavement” [Chen, et al, J. Comput. Phys. 230 (2011) 7018]
  - Particles  $(x, v)$  represented as functions of field  $F(x)$ 
    - DOF reduced to  $N_G$  !!!
- This work: study the **minimum cost** of implicit scheme
  - Focus on Darwin model, instead of Vlasov-Maxwell [Perse, et al, Numkin 2019]
  - Focus on Alfvén modes, kinetic electron, realistic  $\beta/m_e$
  - Derive general equations for FEM-particle-implicit scheme, useful for **TRIMEG** and also **other codes** (Finite difference in [Chen 2011])



Present **TRIMEG** version:  
electrostatic,  $\delta f$ , mixed  
PIC-PIF, unstructured  
mesh, FEM, **as a testbed**.  
[Z.X. Lu, Ph. Lauber, T.  
Hayward-Schneider, A.  
Bottino, M. Hoelzl, *Phys.  
Plasmas*, [26, 122503 \(2019\)](#)]

# Models and equations

- Normalized equations

- Parallel electron motion:

$$\frac{dl}{dt} = v_{\parallel}; l: \text{parallel coordinate}$$

$$\frac{dv_{\parallel}}{dt} = \partial_{\parallel} \delta \phi + \partial_t \delta A_{\parallel}$$

- Ampere's law & quasi-neutrality

$$\nabla_{\perp}^2 \delta A_{\parallel} = C_A \delta J_{\parallel}$$

$$\nabla_{\perp} \cdot \left( \frac{B_0^2}{B^2} \nabla_{\perp} \delta \phi \right) = C_P \delta N$$

- Normalization units

- $R_N = 1m, v_N = \sqrt{2T_e/m_e}$
- $\delta \phi$  normalized to  $m_e v_N^2 / e$

- Crank-Nicolson scheme (implicit)

$$\frac{l^{t+\Delta t} - l^t}{\Delta t} = \frac{v_{\parallel}^{t+\Delta t} + v_{\parallel}^t}{2}$$

$$\frac{v_{\parallel}^{t+\Delta t} - v_{\parallel}^t}{\Delta t} = \partial_{\parallel} \frac{\delta \phi^{t+\Delta t} + \delta \phi^t}{2} + \frac{\delta A_{\parallel}^{t+\Delta t} - \delta A_{\parallel}^t}{\Delta t}$$

$$\nabla_{\perp}^2 \delta A_{\parallel}^{t,t+\Delta t} = C_A \delta J_{\parallel}^{t,t+\Delta t}$$

$$\nabla_{\perp} \cdot \left( \frac{B_0^2}{B^2} \nabla_{\perp} \delta \phi^{t,t+\Delta t} \right) = C_P \delta N^{t,t+\Delta t}$$

- Total Degree of Freedom:  $(N_G + N_p)!$

# Analytic treatment in iteration: “moment enslavement”

## • Iteration steps

$i$ : iteration #

$$\xrightarrow{1} \left\{ \begin{array}{l} \delta\Phi^{start}(t + \Delta t) \\ \delta A_{\parallel}^{start}(t + \Delta t) \end{array} \right\}^i \xrightarrow{2} \left\{ \begin{array}{l} l(t + \Delta t) \\ v_{\parallel}(t + \Delta t) \end{array} \right\}^i \xrightarrow{3} \left\{ \begin{array}{l} \delta N^{end}(t + \Delta t) \\ \delta J^{end}(t + \Delta t) \end{array} \right\}^i \xrightarrow{3} \left\{ \begin{array}{l} \delta\Phi^{end}(t + \Delta t) \\ \delta A_{\parallel}^{end}(t + \Delta t) \end{array} \right\}^i \xrightarrow{4} \left\{ \begin{array}{l} \delta\Phi^{start}(t + \Delta t) \\ \delta A_{\parallel}^{start}(t + \Delta t) \end{array} \right\}^{i+1}$$

1. Each iteration starts with the given field  $\{F^{start}(t + \Delta t)\}^i$ , where  $F \equiv \{\delta\phi, \delta A_{\parallel}\}$
2. Push particles implicitly from  $t$  to  $t + \Delta t$ .
3. In the end of each iteration, fields  $\{F^{end}(t + \Delta t)\}^i$  are solved
4. Fields  $\{F^{start}(t + \Delta t)\}^{i+1}$  are set so that  $|\{F^{start}(t + \Delta t)\}^{i+1} - \{F^{end}(t + \Delta t)\}^{i+1}| \rightarrow 0$  in the next iteration.

Analytic form for setting  $\{F^{start}(t + \Delta t)\}^{i+1}$ :  $\{\delta\phi^{start}, \delta A_{\parallel}^{start}\}^{i+1} = \{\delta\phi^{end}, \delta A_{\parallel}^{end}\}^i + \{\Delta\delta\phi, \Delta\delta A_{\parallel}\}$ , with

$$\left[ \left( \begin{array}{cc} \frac{1}{c_p^2} \nabla_{\perp}^2 & 0 \\ 0 & \frac{1}{c_A} \nabla_{\perp} \cdot \frac{B_0^2}{B^2} \nabla_{\perp} \end{array} \right) - \bar{\bar{M}}_c \right] \cdot \begin{bmatrix} \Delta\delta\phi \\ \Delta\delta A_{\parallel} \end{bmatrix} = \begin{bmatrix} \delta N^{end} - \delta N^{start} \\ \delta J^{end} - \delta J^{start} \end{bmatrix}^i, \bar{\bar{M}}_c = \begin{bmatrix} \frac{\partial \delta N^{end}}{\partial \delta\phi} & \frac{\partial \delta N^{end}}{\partial \delta A_{\parallel}} \\ \frac{\partial \delta J^{end}}{\partial \delta\phi} & \frac{\partial \delta J^{end}}{\partial \delta A_{\parallel}} \end{bmatrix}$$

$(N_G + N_p)^2$  matrix inversion: unacceptable  $\rightarrow$  particle enslavement [G. Chen 2011]

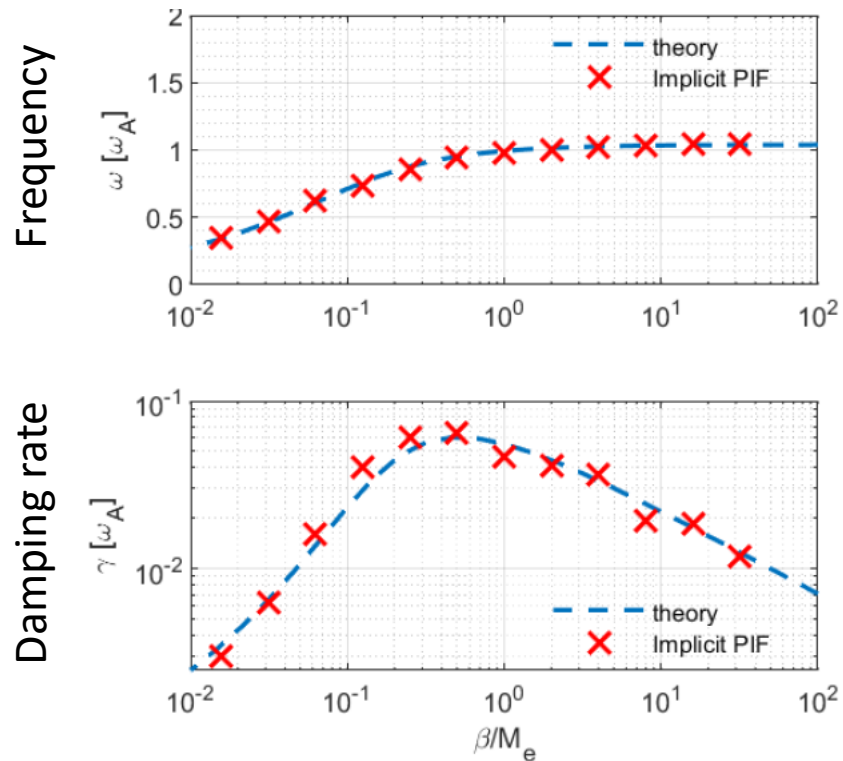
$$\bar{\bar{M}}_c \approx \begin{bmatrix} \frac{k_{\parallel}^2 \Delta t^2}{4} & -i \frac{k_{\parallel} \Delta t}{2} \\ i \frac{k_{\parallel} \Delta t}{2} & 1 \end{bmatrix} \text{ obtained analytically!}$$

$\bar{\bar{M}}_c$  calculation ( $N_G \times N_p$ ): not cheap  $\rightarrow$  analytic acceleration here, i.e., “moment enslavement”

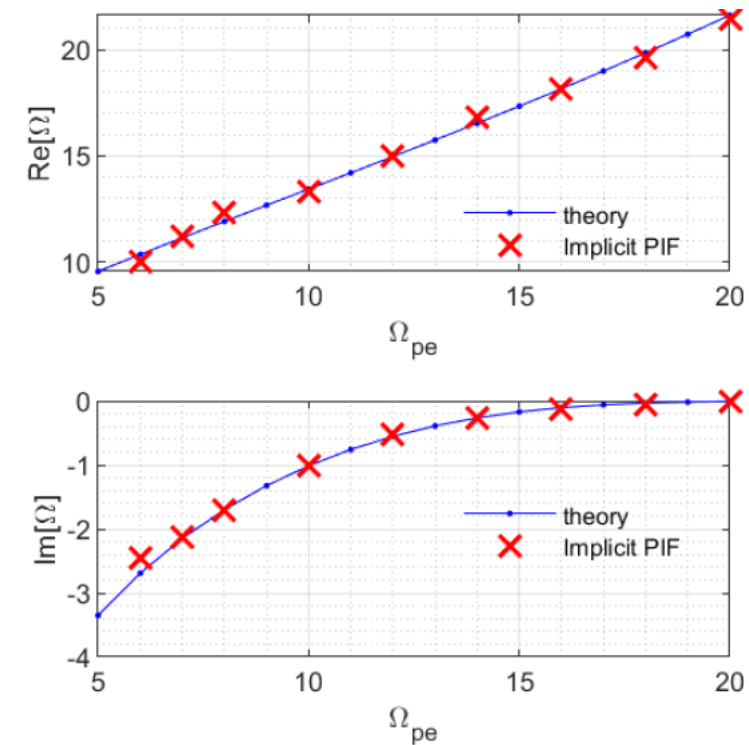
Otherwise more expensive by a factor of  $N_p/N_G$ , i.e., particle/grid number ratio, if calculating  $\bar{\bar{M}}_c$  using particles.

# Shear Alfvén wave simulated

- 1D Shear Alfvén wave simulated
  - Accessible  $\beta/M_e=32$ , i.e.,  $\beta = 0.01$ ,  $M_e \equiv \frac{m_e}{m_i} = 1/3200$



Shear Alfvén wave

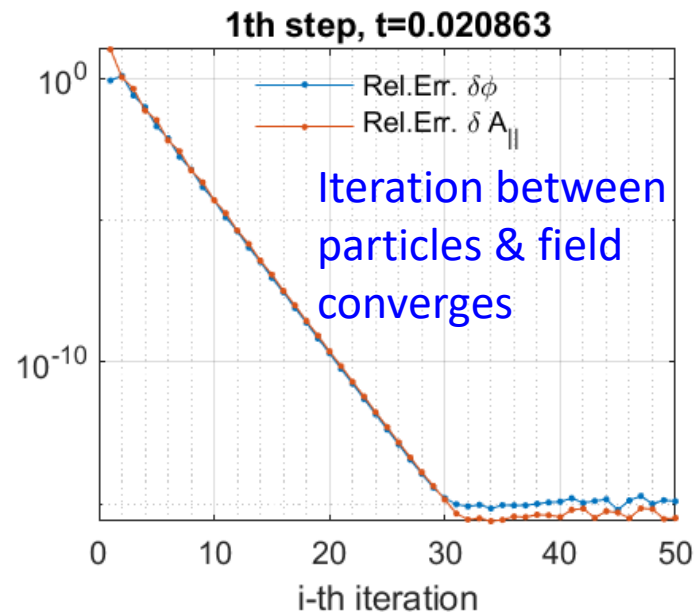


Electron Langmuir wave

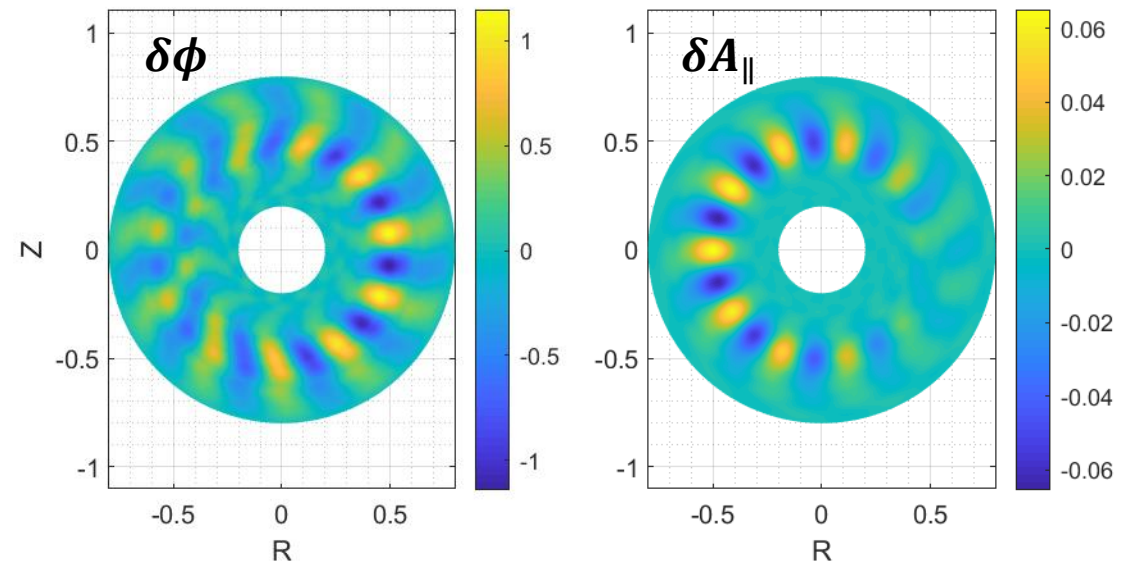
Non uniform loading needed for resolving wave-particle resonance

# Shear Alfvén wave simulation (prototype version)

- Simplified model for Toroidicity induced Alfvén Eigenmode (TAE) simulation
  - Ion response: only polarization density
  - Electron: full f, parallel nonlinear motion kept (zero drift & ExB motion; zero mirror force)
- Parameters: aims for ITPA case [Koenies et al 58, 126027(2018)]
- Convergence of the implicit scheme
- TAE 2D mode structure (w/o EPs)



$$Rel. Err. \{\delta\phi\} = \frac{\sqrt{\int |\Delta\delta\phi|^2 dS}}{\sqrt{\int |\delta\phi|^2 dS}}$$



Caveat: reduced  $\beta/m_e$  value with  $m_e = 0.002$ ,  $\beta = 0.001$ .  
Fourier in  $(\theta, \phi)$ , FEM in  $r$  direction. Toroidal mode #:  $n = -6$ .  
Only 2 poloidal harmonics kept.

# Mixed explicit (Runge-Kutta)-implicit scheme in torus

- Treat the equilibrium particle motion explicitly but the perturbed one implicitly

$$\frac{d\mathbf{R}}{dt} = v_{\parallel} \mathbf{b} + \mathbf{v}_d + \frac{\mathbf{b}}{B} \times \nabla(\delta\phi - v_{\parallel} \delta A)$$

$$\frac{dv_{\parallel}}{dt} = -\mu \partial_{\parallel} B - \frac{e_s}{m_s} (\partial_{\parallel} \delta\phi + \partial_t \delta A)$$

It is possible to treat all terms implicitly, but not necessary

- Equilibrium particle motion (**explicit**)

$$\frac{d\mathbf{R}_0}{dt} = v_{\parallel} \mathbf{b} + \mathbf{v}_d$$

$$\frac{dv_{\parallel 0}}{dt} = -\mu \partial_{\parallel} B$$

- Perturbed particle motion (**implicit**)

$$\frac{d\mathbf{R}_1}{dt} = \frac{\mathbf{b}}{B} \times \nabla(\delta\phi - v_{\parallel} \delta A)$$

$$\frac{dv_{\parallel 1}}{dt} = -\frac{e_s}{m_s} (\partial_{\parallel} \delta\phi + \partial_t \delta A)$$

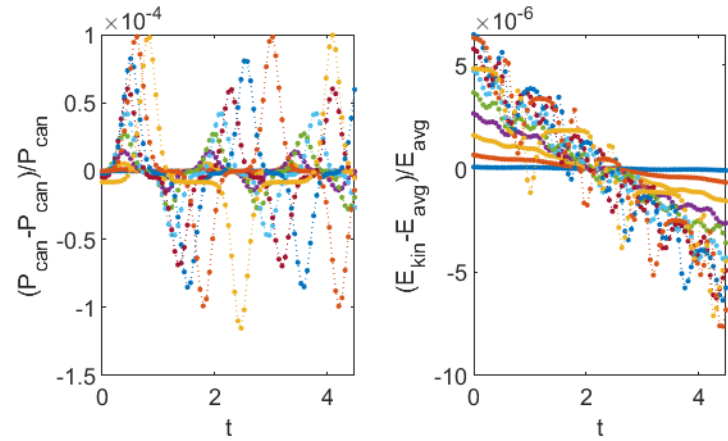
Runge-Kutta 4<sup>th</sup> order implemented  
Geometric scheme of interest



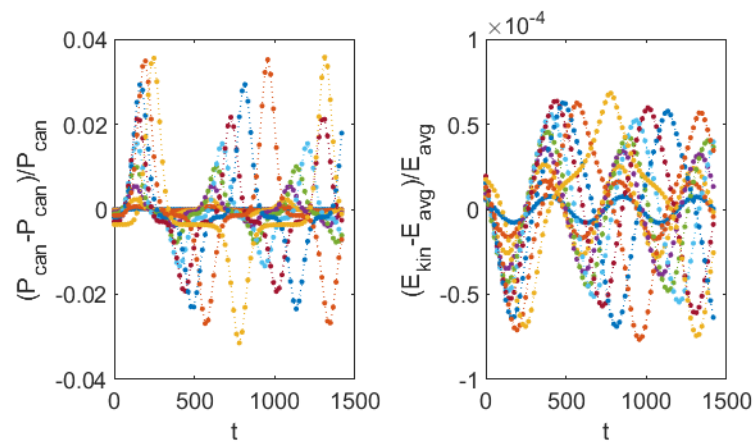
# RK4 particle integrator, implicit particle-field iteration

- RK4 integrator gives **reasonable conservation** of canonical momentum ( $P_\phi$ ) & energy ( $E$ )

- The implicit field-particle solver shows reasonable convergence
  - Radial component of  $\delta v$  included
  - More tests for **large  $\beta/m_e$**  ongoing

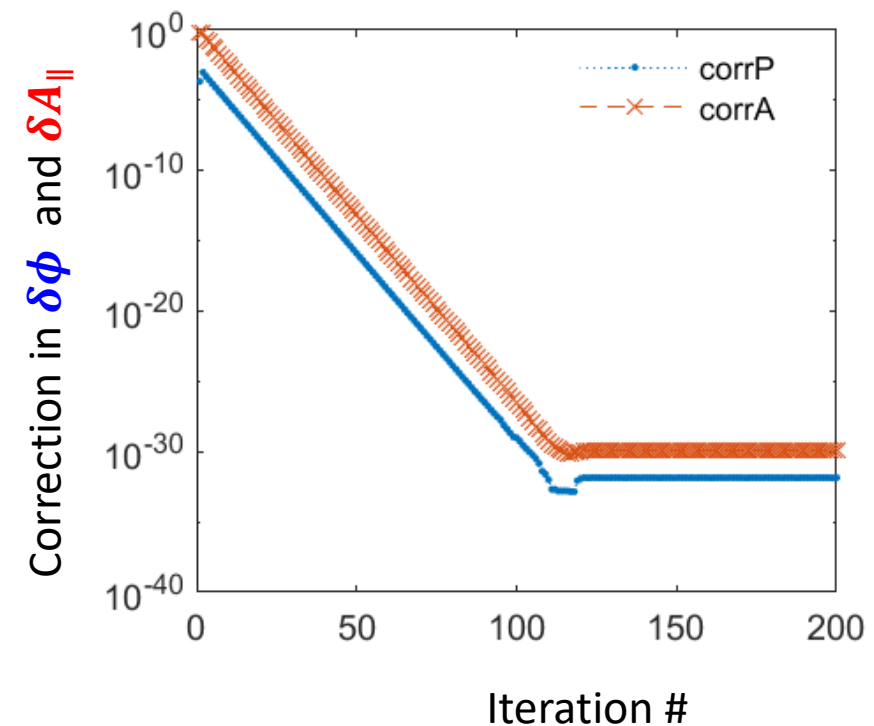


$$\rho_t = 6.3e-5$$



$$\rho_t = 0.02$$

$\rho/a$  terms needed for larger  $\rho_t$

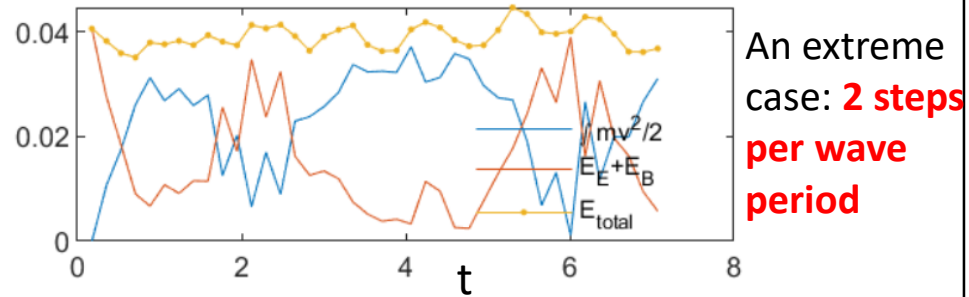


# Outlook

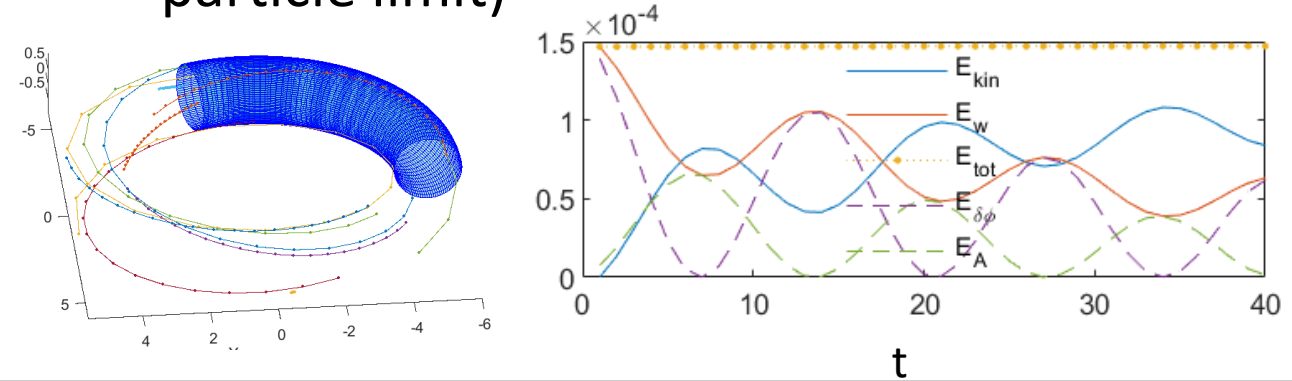
- Possible contributions to MET in 2020
  - Full  $f$ , implicit, nonlinear particle simulation of Alfvén waves, driven by EPs
    - Caveat: single  $n$  (particle-in-Fourier in toroidal direction); ad-hoc equilibrium
    - In parallel with HAGIS full  $f$  simulation
  - Anisotropic EP distribution effects on zonal flow residual
    - Previous IAEA work to be finalized [Z. Lu et al, Effects of anisotropic energetic particles on residual zonal flow]
- Comments & suggestions are welcome

# Four tests of implicit scheme (full f)

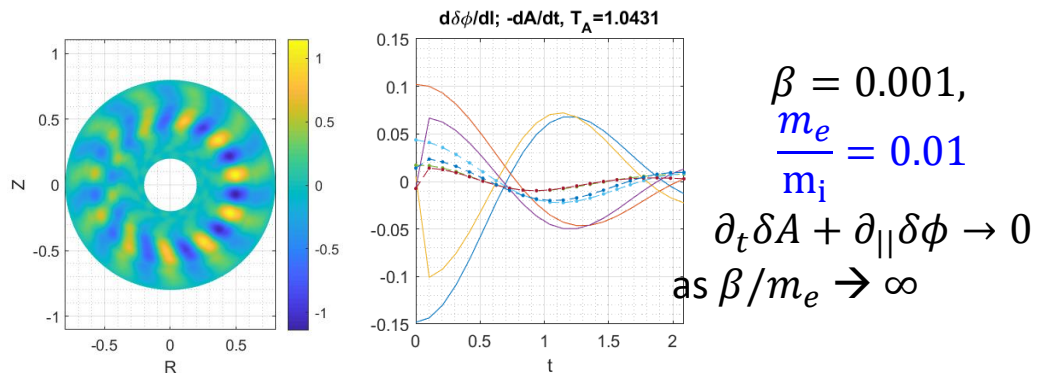
- 1D case, gyrokinetic electrons
  - Shear Alfvén wave, SAW (done)
  - Current driven mode (to do)



- 2D Single flux surface (cylinder limit)
  - Shear Alfvén wave simulated (well passing particle limit)



- 3D circular geometry (tokamak)
  - TAE structure, simplified model
  - Need improvement for large  $\beta/m_e$



- 3D realistic tokamak
  - Need more time in 2020
  - Goal: tests of electromagnetic, full f, implicit, kinetic  $e^-$  model
  - Complementary to other codes

