

# 3D nonlinear MHD studies at Consorzio RFX: Achievements and challenges in macroscopic modelling of fusion plasmas

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## ABSTRACT

**Macroscopic helical self-organization** occur in pinch configurations: Eminently in Reversed Field Pinches, and in Tokamak scenarios too, for example when dynamo/flux-pumping effects play a role. Main achievements in 3D nonlinear MHD have been obtained within **visco-resistive full-MHD (SpeCyl and PIXIE3D codes)** [2,3]

- Transition to Quasi Helical regimes in RFP (QSH)[6,7];
- Boundary Conditions extension toward realistic RFX-mod2 front-end (RFP and Tokamak) [10,11];
- Formation of Internal Transport Barriers in RFP (eITB), (Lagrangian Coherent Structures)[4,8];
- Large scale modes (dynamo/flux pumping effect), their control and characterization of plasma flow: Magnetic Reconnection: RFPs and tokamaks, interplay with Alfvén waves (possible RFP ion heating mechanism) [9];

Then, significant steps forward have been achieved within reduced-MHD **JOREK** code applications, important in view of DTT:

- Assessment of **Shattered Pellet Injection** to mitigate disruptions, and comparison vs the MGI approach (extending JET applications [5,12]);
- Initial assessment about **E.L.M physics and correlation with 3D fields**. [13]

These lines of research are naturally interconnected with nonlinear MHD activities under EUROfusion (TSVV, WPTE) programmes.

References: [1] 2023 RFP Boundary Physics and MHD, disruptions and control physics; [2] 2023 RFP Boundary Physics and MHD, disruptions and control physics; [3] 2023 RFP Boundary Physics and MHD, disruptions and control physics; [4] 2023 RFP Boundary Physics and MHD, disruptions and control physics; [5] 2023 RFP Boundary Physics and MHD, disruptions and control physics; [6] 2023 RFP Boundary Physics and MHD, disruptions and control physics; [7] 2023 RFP Boundary Physics and MHD, disruptions and control physics; [8] 2023 RFP Boundary Physics and MHD, disruptions and control physics; [9] 2023 RFP Boundary Physics and MHD, disruptions and control physics; [10] 2023 RFP Boundary Physics and MHD, disruptions and control physics; [11] 2023 RFP Boundary Physics and MHD, disruptions and control physics; [12] 2023 RFP Boundary Physics and MHD, disruptions and control physics; [13] 2023 RFP Boundary Physics and MHD, disruptions and control physics.

## Some open issues, Ongoing works:

**Visco-Resistive 3D nonlinear MHD** provides a fruitful global description of relaxation-reconnection features consistent with several aspects typical of both RFP and Tokamak configurations. Several aspects await for **refined and extended modelling**:

- Upgrade toward **realistic Boundary Conditions**
- Verification SpeCyl – PIXIE3D (Chacon-LosAlamos)
- proof of principle studies on free-boundary modes Spinicci, Bonfiglio et al. *AIP Advances* (2023), *J. Plasma Phys.* (2024)
- Further develop **analogies RFP – Tokamak**
- Alfvén waves excitation at sawtoothing (ohmic) Kryzhanovskyy, Bonfiglio et al. *Nuclear Fusion* 2022, Kryzhanovskyy, Bonfiglio et al. *Nuclear Fusion* 2024
- Assess **magnetic topology characterization (and related transport) - Lagrangian Coherent Structures (LCS) validation on experimental eITB** follow up of Veranda et al. *NF* 2017
- Extended modeling to be addressed with the 3D codes: PIXIE3D and, possibly, JOREK

Address **momentum transport impact and estimate experimental effective Hartmann  $H = (\nu \eta)^{-1/2}$**

Momentum transport in plasmas is a longstanding open issue. Heuristic approach proposed in:

- Terranova, Bolzonella, Cappello et al. *PPCF* 2000
- Scaling of magnetic fluctuations in the RFX reversed field pinch - Vivenci, Spizzo, Veranda et al. *Theory Fusion Plasmas* 2022 *JPCS*
- Vivenci, Veranda, Cappello Bonfiglio PoP 2023

What is the mechanism of “anomalous” ion heating in RFPs (observed at sawtoothing)?

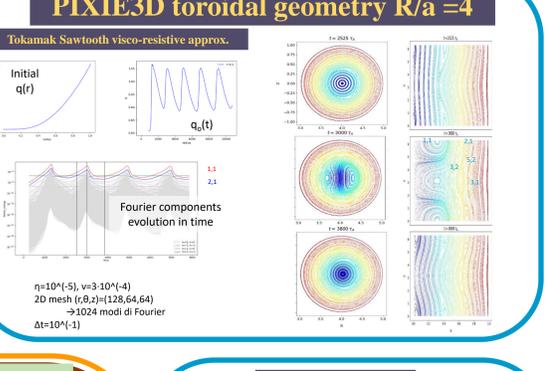
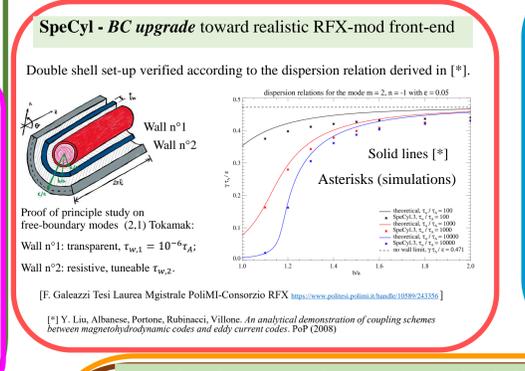
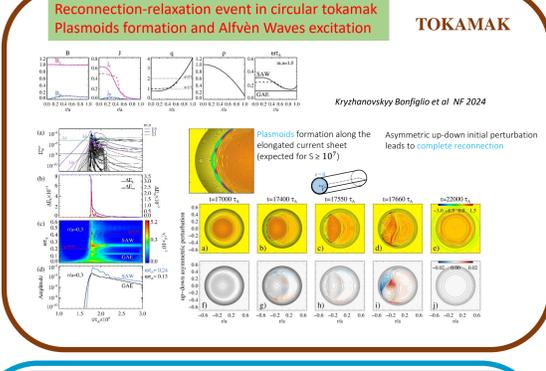
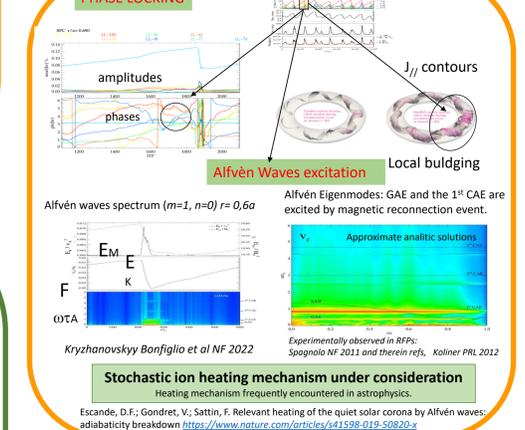
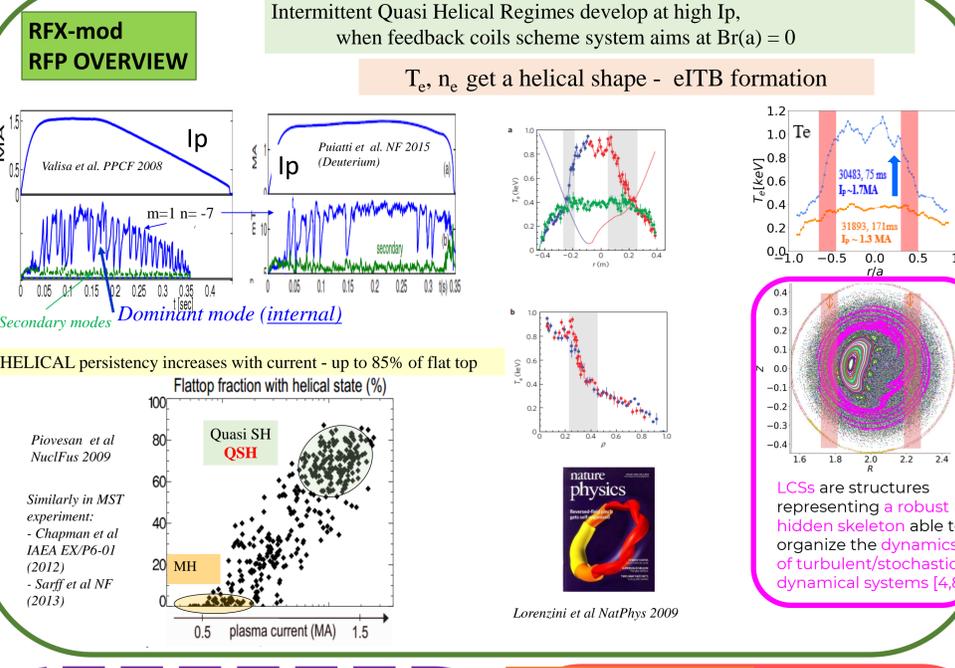
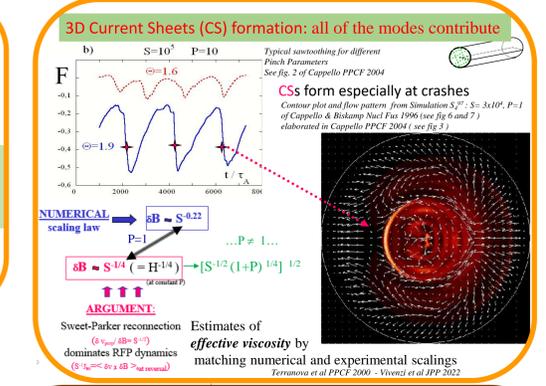
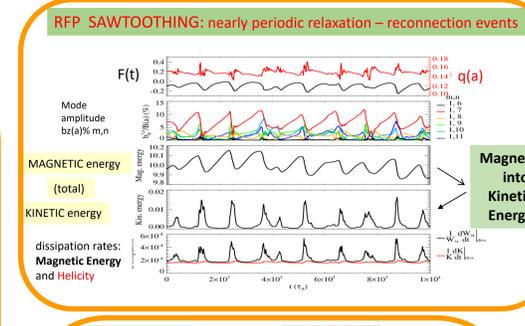
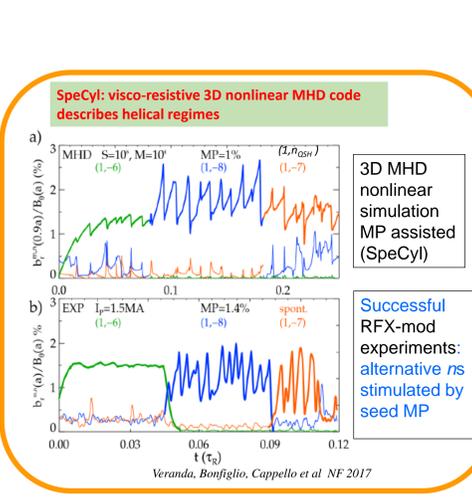
Non resonant Alfvén wave particle interaction Veranda, Sattin et al. *in preparation*

**THIS CONFERENCE**  
L Carraro Spettroscopia come finestra sul plasma: diagnosi e trasporto delle impurezze 5 febbraio 10.40

### INTRODUCTION

**RFP experiments:** at high current ↔ ordered KINKED plasma advanced operation with **Boundary Control**

**RFX-mod:** CLEAN MODE CONTROL and/or NON CONVENTIONAL SCENARIOS (PPCD-OPCD)



### Toy model: the natural kinking of the current carrying plasmas in a flux conserver

- 1) Kink unstable
- 2) solenoidal effect by the wire itself
- 3) Kink saturates if edge field reversal is forced

for too small  $\frac{L}{a}$  → no reversal  $\rightarrow a$  (disruption)

... Tokamak case ...

### 3D nonlinear MHD

SpeCyl code visco-resistive approx. Cappello & Biskamp *Nucl. Fus.* 1996

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \wedge (\mathbf{v} \wedge \mathbf{B}) - \nabla \wedge (\boldsymbol{\eta} \mathbf{J})$$

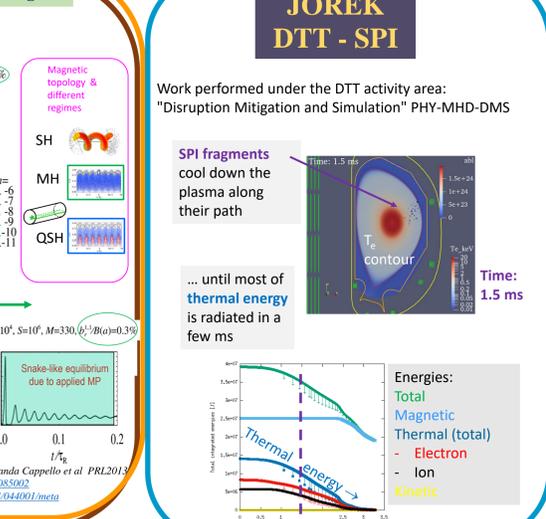
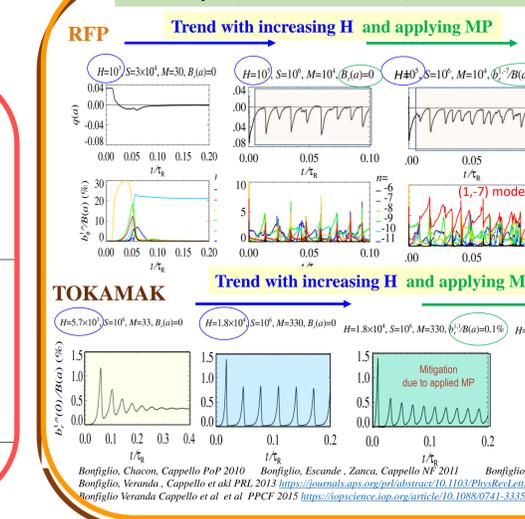
$$\frac{d\mathbf{v}}{dt} = \mathbf{J} \wedge \mathbf{B} + \nabla^2 \mathbf{v}$$

$\rho \equiv 1, \nabla \cdot \boldsymbol{\rho} = 0$

“typical” boundary conditions:

- $B^z = 0$  (constant magnetic flux  $\Phi$ )
- Constant  $E_z$
- Ideal boundary
- MP on  $Br_{min}$  (~ 1%, 2%, 4 %)
- Thin shell + vacuum + ideal wall
- velocity field: no slip.

initial conditions decide  $\Phi, I_z$



After kinking ...

“slinky” instability

localized collapse

bulging-collapse at toroidal location, similar in experiments and nonlinear modeling...

Some additional references from ECMP 2025 ->

### The Hartmann number

$$\mathbf{t} \rightarrow \tilde{\mathbf{t}} = \frac{\eta}{\nu} \mathbf{t}, \quad \mathbf{v} \rightarrow \tilde{\mathbf{v}} = \frac{\nu}{v} \mathbf{v}$$

$(S, M) \rightarrow (\eta, \nu) \rightarrow (H, P)$

Prandtl\_mag  $P = \nu \eta / S = M$

Hartmann  $H = (\nu \eta)^{-1/2}$

H regulates both current density and momentum diffusion

H measures the relative weight of:

- electromagnetic Lorentz forces
- visco-resistive friction
- the balance between non-linear mode coupling and - damping / smoothing yielding the actual spectral broadening of MHD activity, excited by linear instability or external means (RMP)

$$\frac{1}{P} \frac{d\tilde{\mathbf{v}}}{dt} = \tilde{\mathbf{J}} \wedge \tilde{\mathbf{B}} + \nabla^2 (\tilde{H}^2 \tilde{\mathbf{v}})$$

$\rho \equiv 1, \nabla \cdot \boldsymbol{\rho} = 0$

H rules the amplitude of the nonlinearly driven  $m=0$  modes

