

# ASTROPHYSICAL JET MODELING

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1. **What** we mean by *jet* ?

Astrohysical jets are collimated supersonic (relativistic) flows along which mass, momentum, energy and magnetic flux are channeled from several astrophysical objects to outer medium.

2. **Where** do we find *jets* ?

Astrohysical jets are present in many astronomical objects, they span a wide range of spatial scales: on stellar scale we find them in the solar corona, in the first phases of stellar evolution (YSO) or in the last phases of stellar evolution (X-ray binaries, X-ray transients, Pulsars, GRBs ...). On the extragalactic scale (AGN jets)

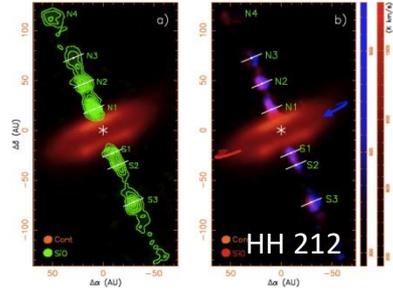
3. **Why** do we study *jets* ?

On one side they can provide information about their sources, on the other side they may have a strong impact on the surrounding medium.

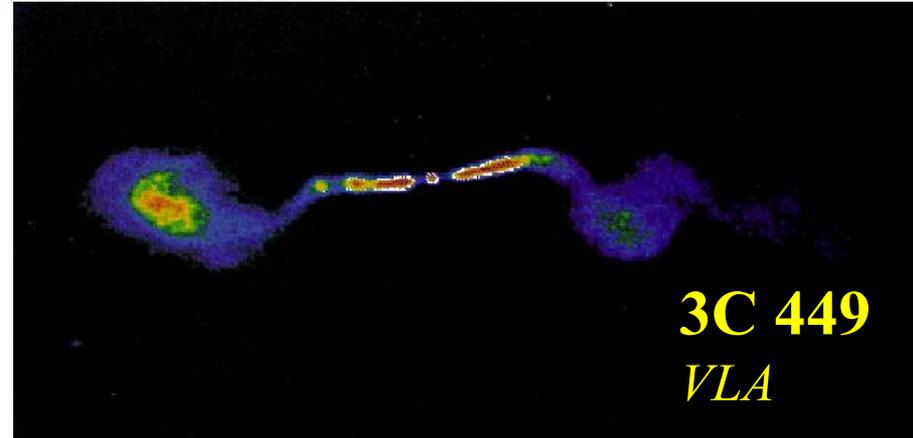
## Stellar jets



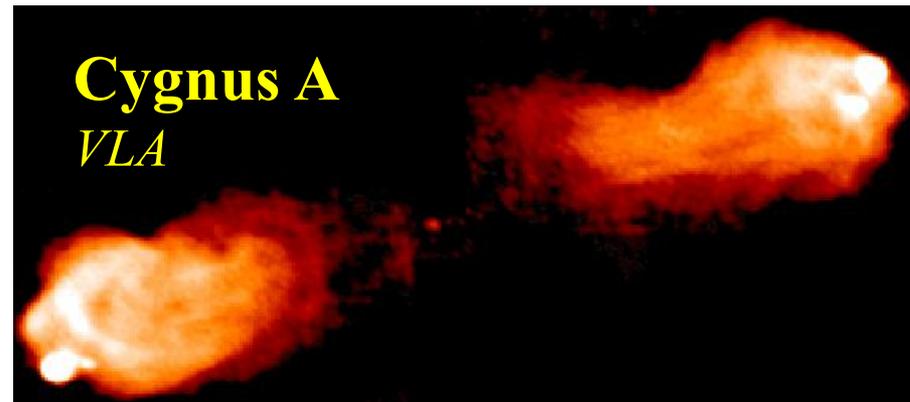
Ray & Ferreira, 2021



## AGN jets



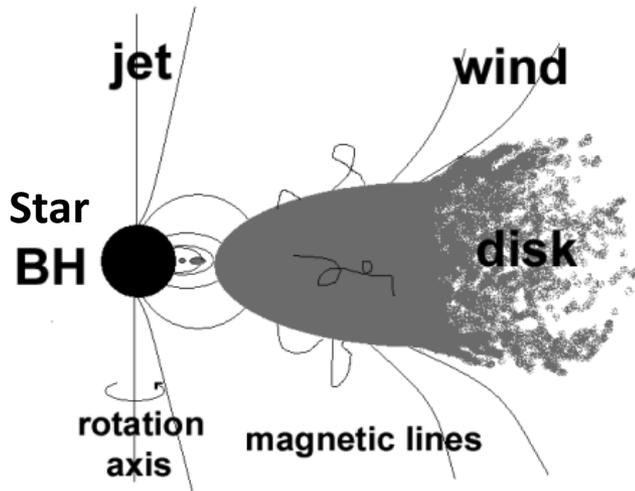
X-ray: NASA/CXC/SAO  
Optical: NASA/STSCI  
IR: NASA/JPL/Caltech



## Similarities and differences

	YSO jet $\leq$ pc	AGN jets $>$ kpc
Central object	Star	SMBH
Radiation	Thermal	Non-Thermal
Spectrum	Optical, IR	Radio to gamma rays
Compositions	(e-,p+)	(e-,p+), (e-,e+), e.m.
Jet power	$10^{30-33}$ erg/s	$10^{40-47}$ erg/s
Mach Number	20-40	$\gg 1$ (relativistic)
Density ratio (jet/am)	ca 1	$\ll 1$

## The jet-disk paradigm

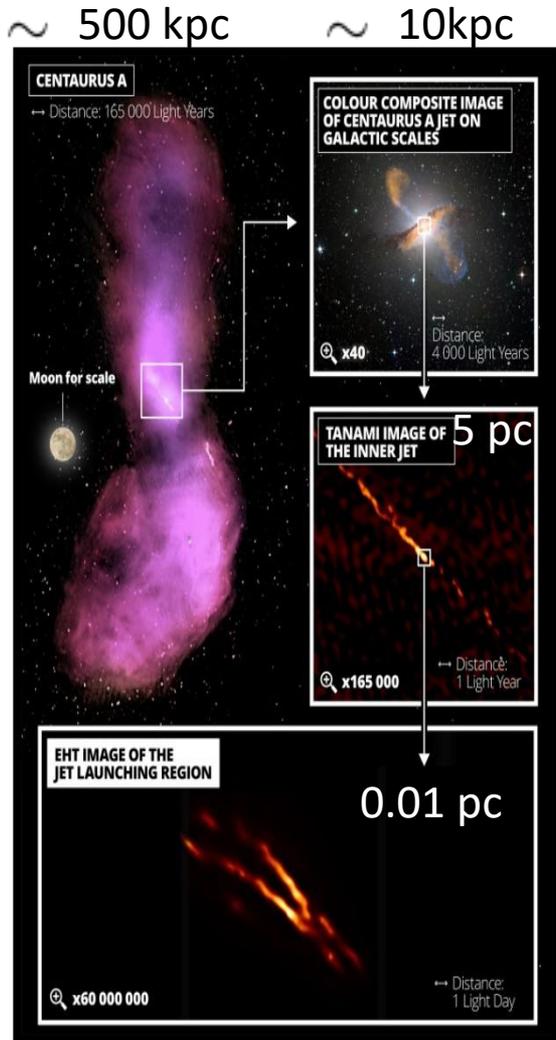


Ingredients:

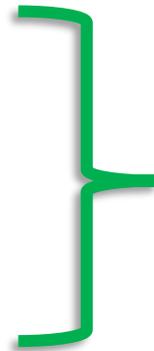
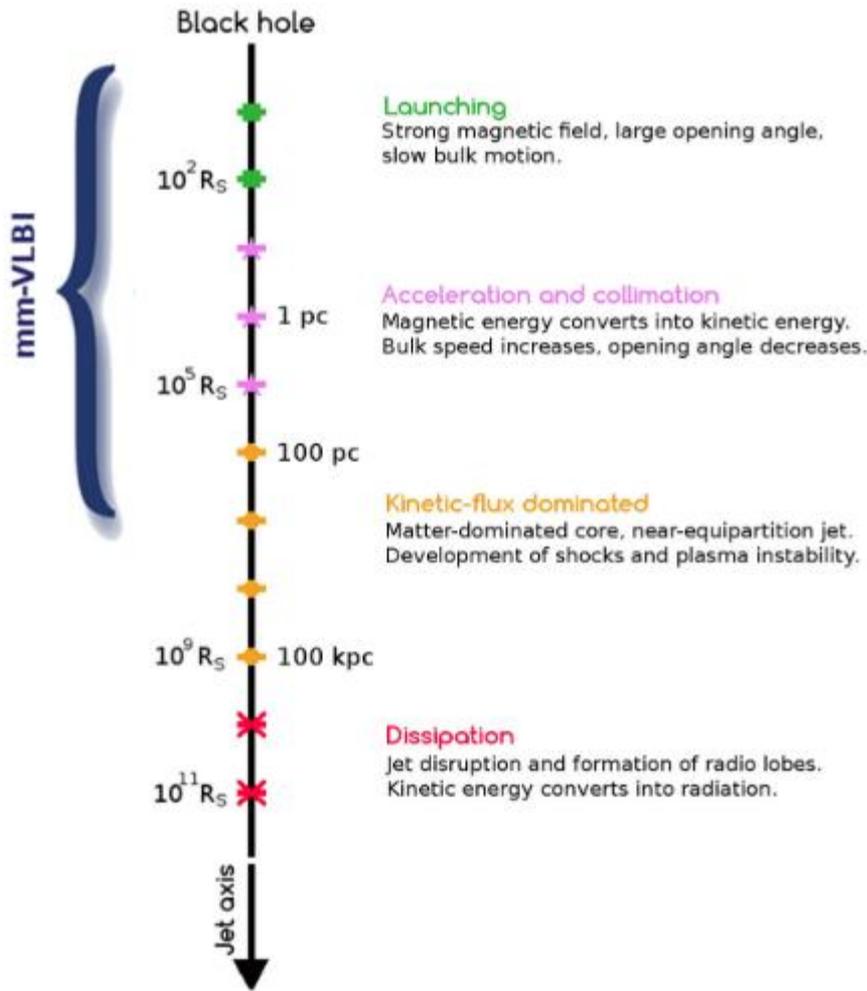
- Central object creating a strong gravitational field
- Accretion disk
- Magnetic field (turbulent in the disk, ordered in magnetosphere)
- Strongly collimated outflow along the rotation axis
- Collimated wind arising from the disk

**Accretion and ejection are strongly connected**

# How can we approach the jet physics?



- **Launching phase:** acceleration from BH/disk and collimation  $\sim 10^{15}$  cm
- **Propagation phase:** confinement and stability  $\sim 1-100$  kpc
- **Termination:** Termination interaction with external medium  $> 100$  kpc ( $3 \times 10^{23}$  cm)
- **High-energy electron acceleration**  $\sim 10^4$  cm



### Launching phase:

Acceleration either from disk or BH  
And collimation

### Propagation phase:

Confinement and stability

### Termination:

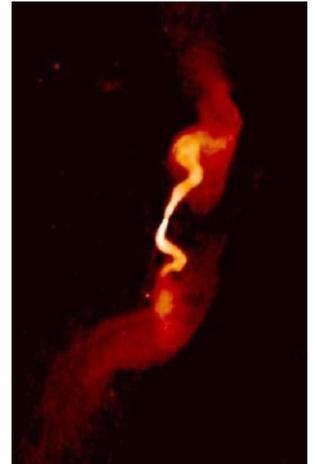
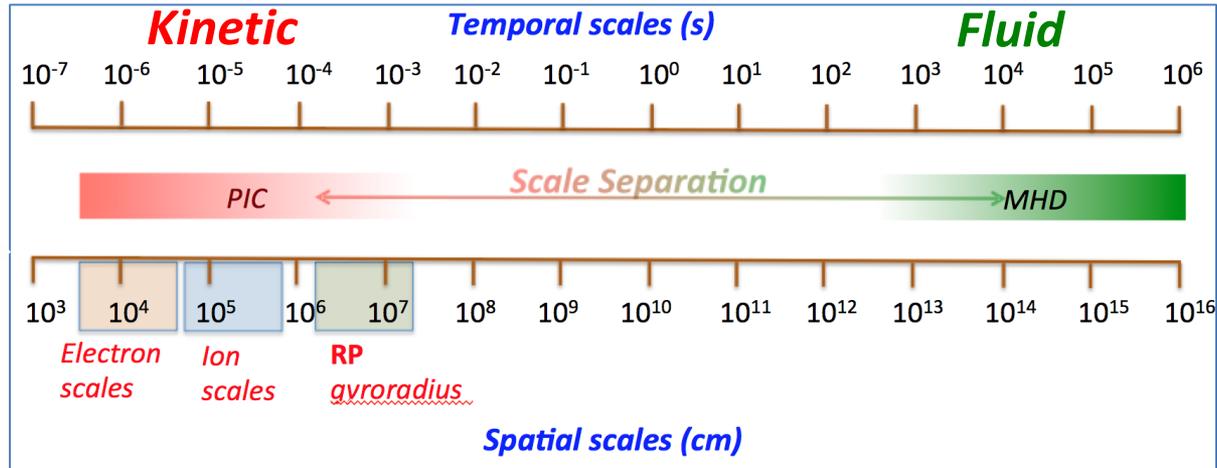
Interaction with external medium

### Non thermal emission:

High energy electron acceleration

Modeling through relativistic MHD simulation

Astrophysical environments involve physical processes operating at *extremely different spatial* and *temporal scales*, and complex *interactions* between *plasmas* and *radiation*.



10<sup>1</sup> → 10<sup>2</sup> s  
10<sup>4</sup> → 10<sup>7</sup> cm

← How we can connect? →

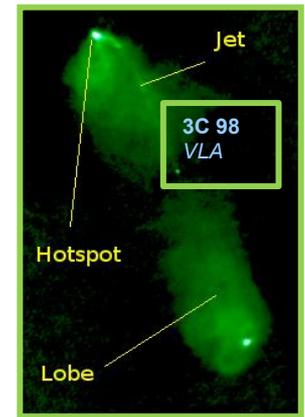
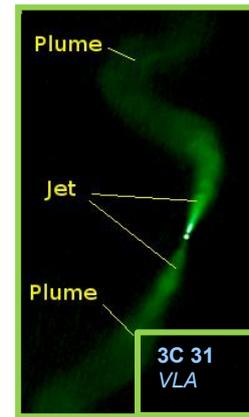
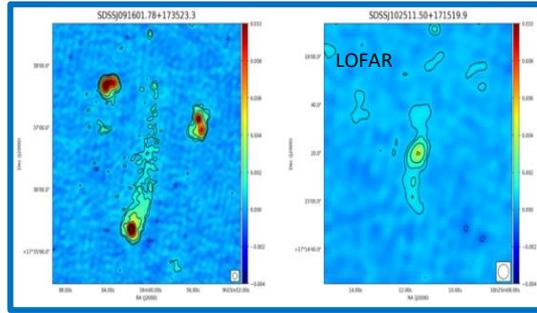
10<sup>10</sup> → 10<sup>15</sup> s  
10<sup>18</sup> → 10<sup>24</sup> cm

Particles acceleration  
And radiation processes

Ideal MHD modeling  
No resistivity, no viscosity  
PLUTO code

# Warning!!!

Also the macrophysics involves a huge range of scale

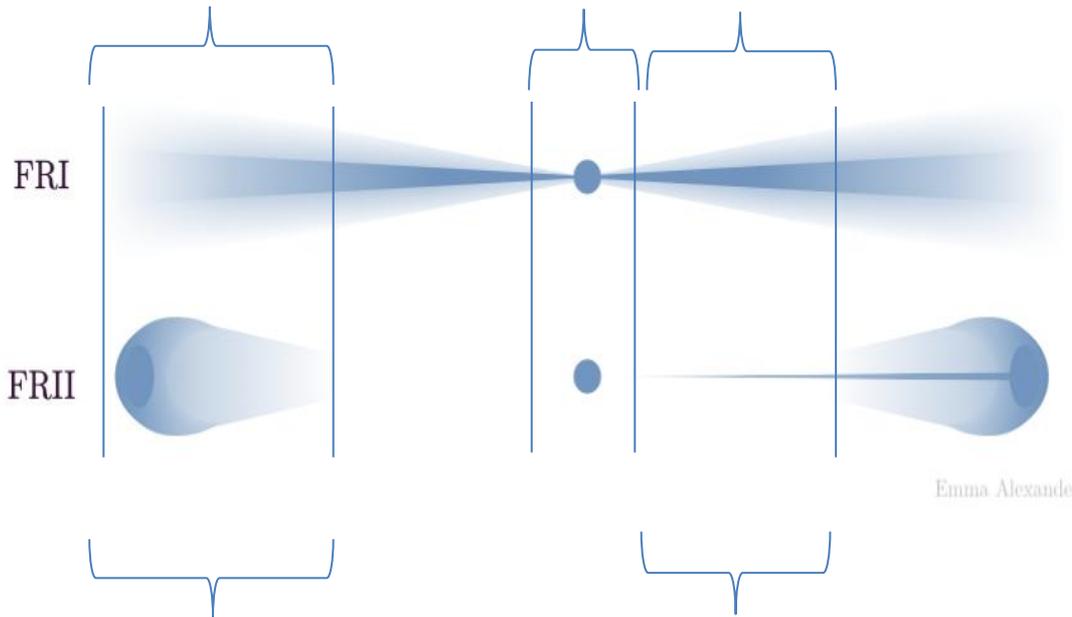


## Multiscale approach

Turbulent propagation up to 100 kpc

Launching region  $10^{15}$  cm

Propagation region with deceleration  $< 1$  kpc



Termination shock, hot spot

Propagation region with No Deceleration 1-100 kpc

At the base all jets start relativistic with similar Lorentz factor (Giovannini+, 2001)

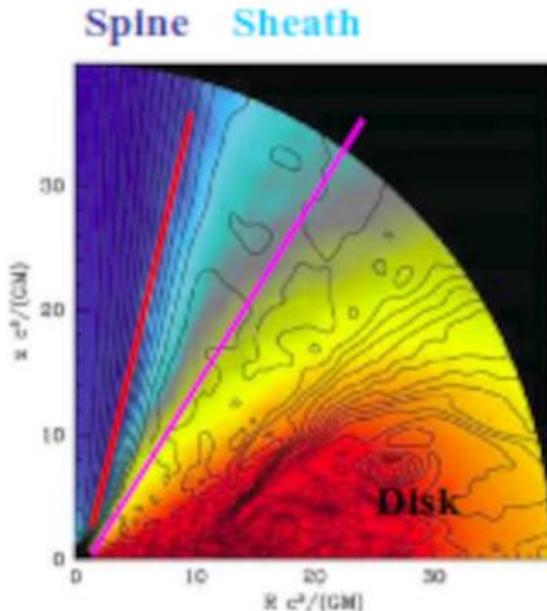
# Launching phase

Main ingredients:

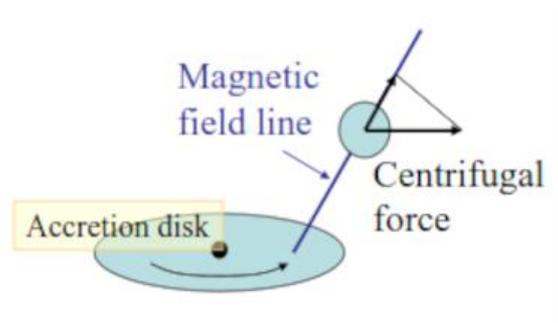
- Rotation
- Magnetic field

Origin:

- BH (Blandford & Znajek 1977) → extraction of BH rotational energy
- Disk (Blandford & Payne 1982) → magnetocentrifugal acceleration



Density distribution  
(McKinney 2006)



Many GRMHD simulations of jet formation (Port+,2019...) suggest that :  
jet spine (Poynting flux jet) is driven by BZ process may be surrounded by a broad sheath wind driven by the magnetic field anchored in the accretion disk.

Open problem:

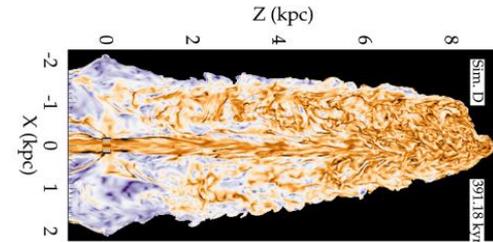
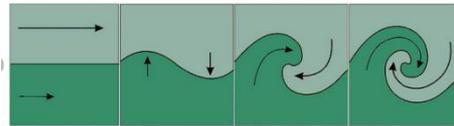
**How the Poynting flux turn in a matter dominated jet?**

Possibility: CD instability due to the toroidal magnetic field

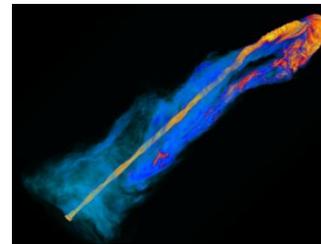
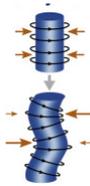
# Propagation phase

The MHD ins. are very important in determining morphology, dynamics and emission properties of astrophysical jets.

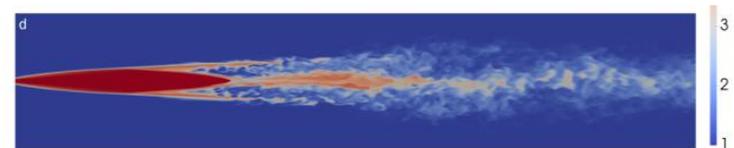
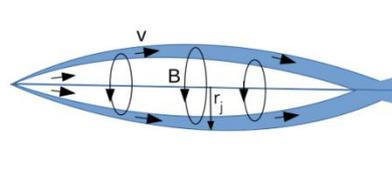
## KELVIN HELMHOLTZ

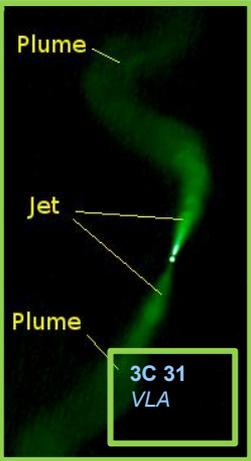


## CURRENT DRIVEN

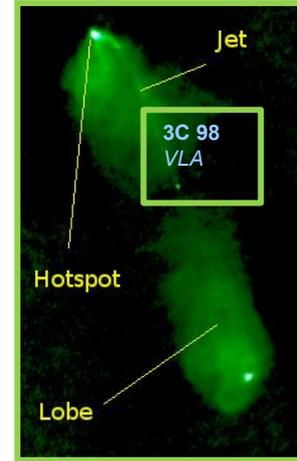


## CENTRIFUGAL

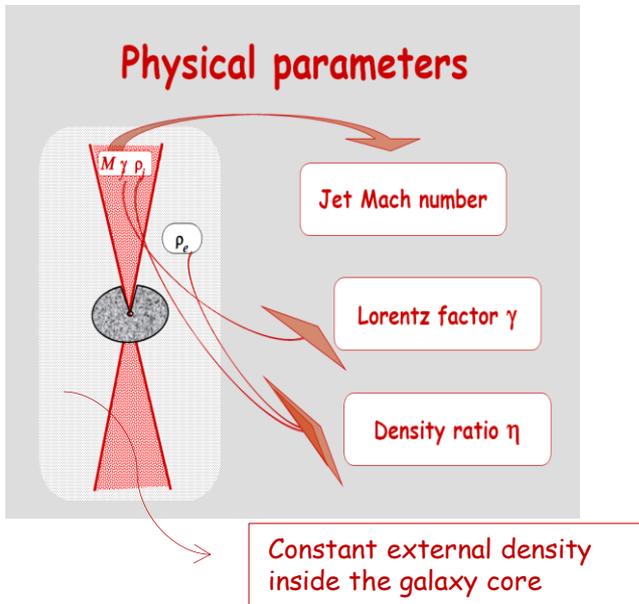
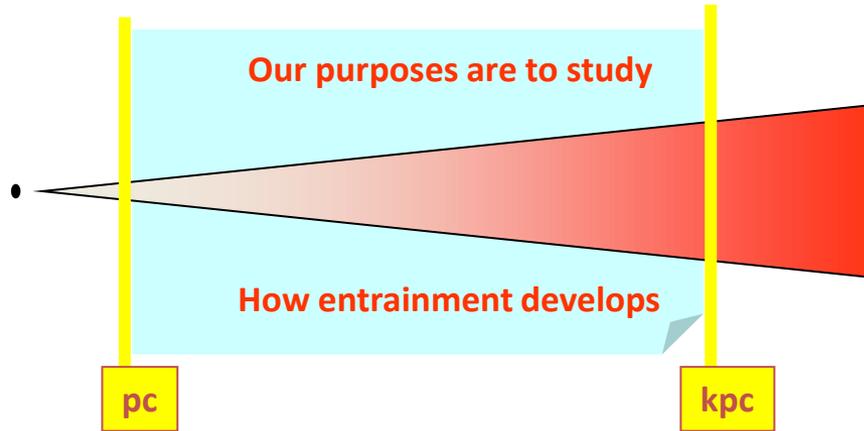




Is it possible to decelerate a relativistic jet to sub-relativistic speed in the first kpc?



### Deceleration



$$P_j \sim h \left( \frac{\gamma}{10} \right)^2 \left( \frac{\eta}{10^{-2}, 10^{-4}} \right) \left( \frac{r_j}{1 \text{ pc}} \right)^2 \left( \frac{n_0}{1 \text{ cm}^{-3}} \right) 1.22 \times 10^{43} \text{ erg s}^{-1}$$

$\downarrow$                        $\downarrow$                        $\rightarrow$   $10^{43}, 10^{45} \text{ erg s}^{-1}$

**Jet instabilities:** growth length-scale increases with Lorentz factor and Mach number



Non linear growth



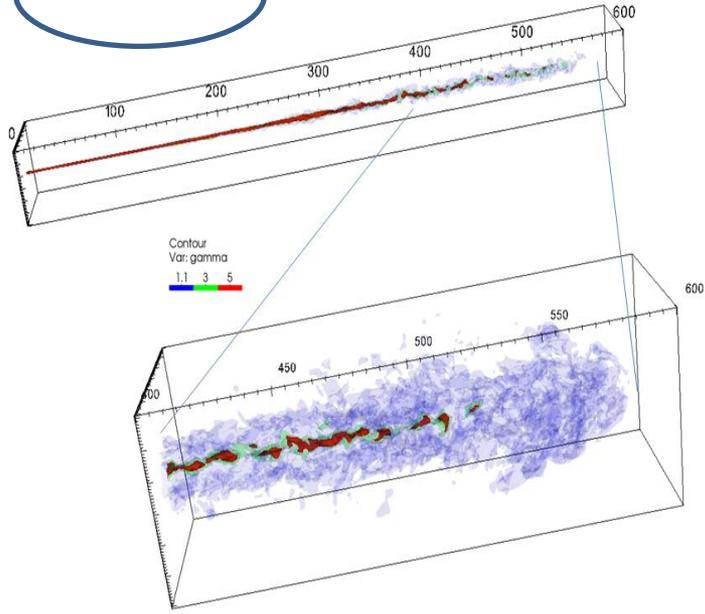
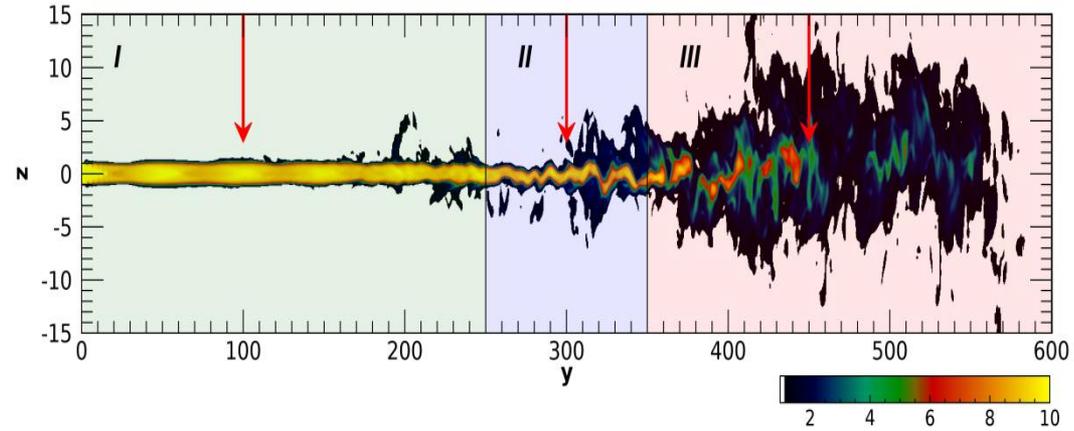
Mixing and mass entrainment



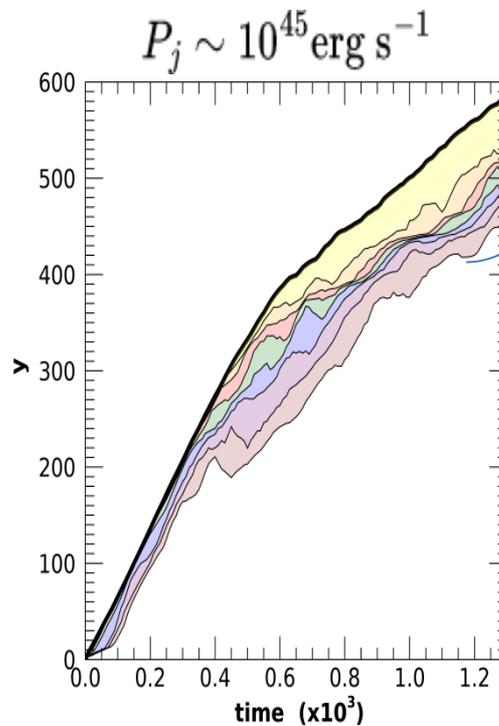
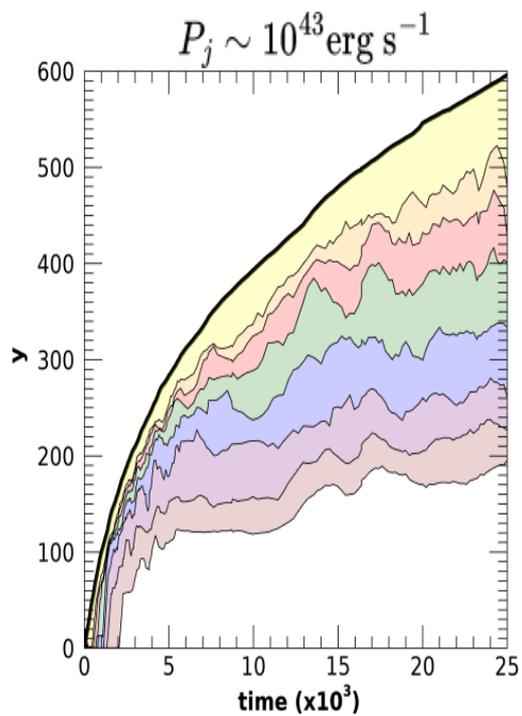
Jet deceleration

Lorentz factor

$$P_j \sim 10^{43} \text{ erg s}^{-1}$$



- I ricollimation shocks**
- II growth of instabilities**
- III mixing and jet breaking**

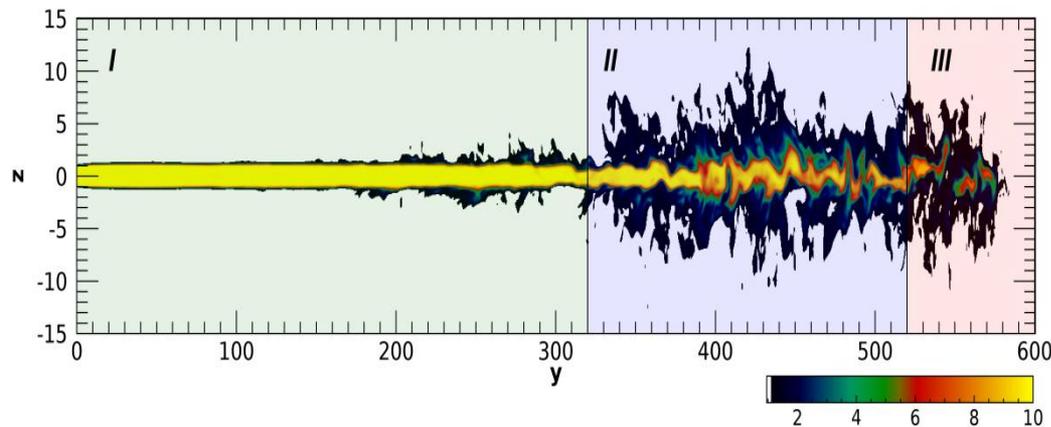


Each curve shows, as a function of time, the distance up to which relativistic material ( $\gamma > 5$ ) carries more than certain fraction of the initial momentum flux: 70%, 60%, 50%.....10%

$P_j \sim 10^{43} \text{ erg s}^{-1}$

- Strong deceleration
- Strong entrainment
- Steady state

Time unit:  $\tau \sim 3.25 \left( \frac{r_j}{1 \text{ pc}} \right) \text{ yr}$



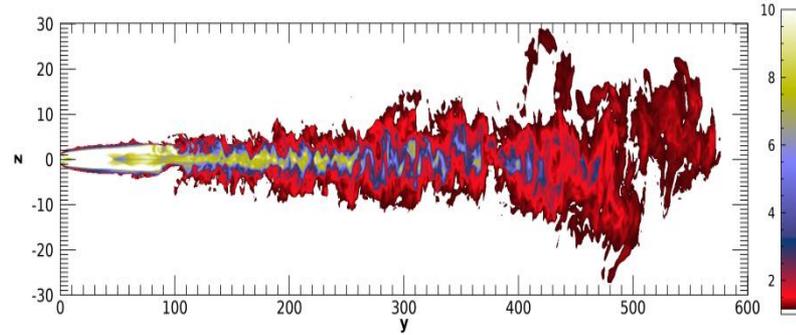
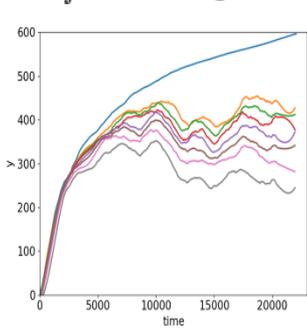
$P_j \sim 10^{45} \text{ erg s}^{-1}$

- Weak deceleration
- Weak entrainment
- No Steady state

# Magnetic effects

$$P_j \sim 10^{43} \text{ erg s}^{-1}$$

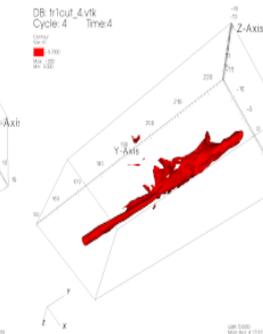
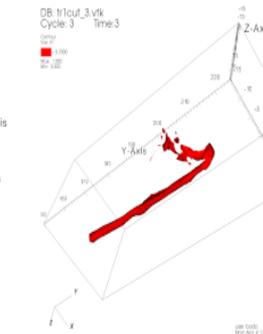
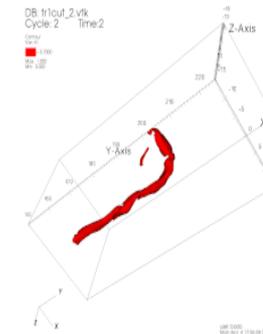
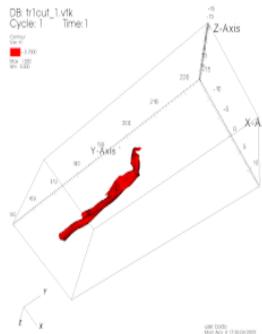
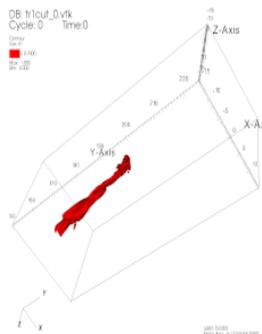
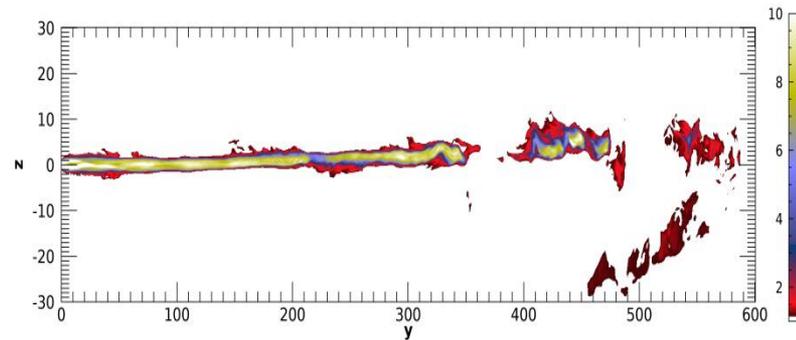
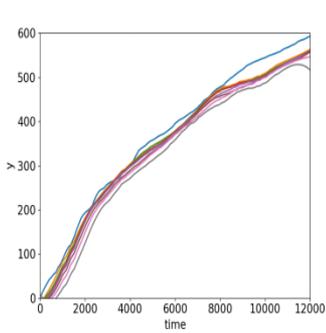
$$\sigma = 10^{-2}$$



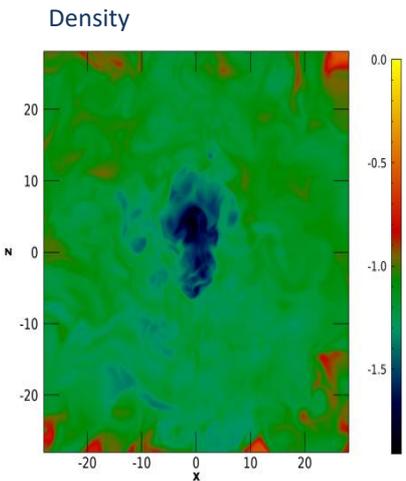
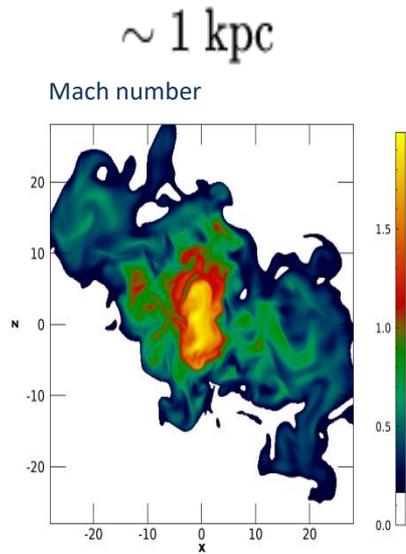
Magnetic field reduces entrainment, stabilizing KH inst. .  
New inst. (current-driven) starts deforming the jet without deceleration.

$$P_j \sim 10^{43} \text{ erg s}^{-1}$$

$$\sigma = 10^{-1}$$



# Turbulent propagation starting from 1 kpc



Initial conditions:

$$M \sim 2$$

$$\eta \sim 10^{-2}$$

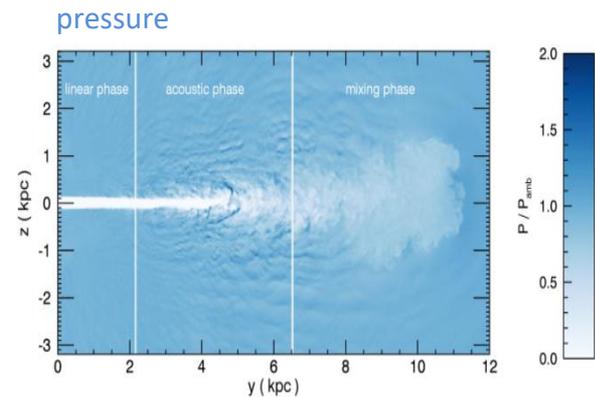
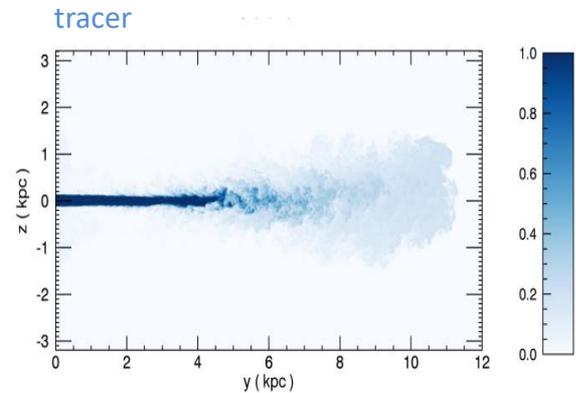
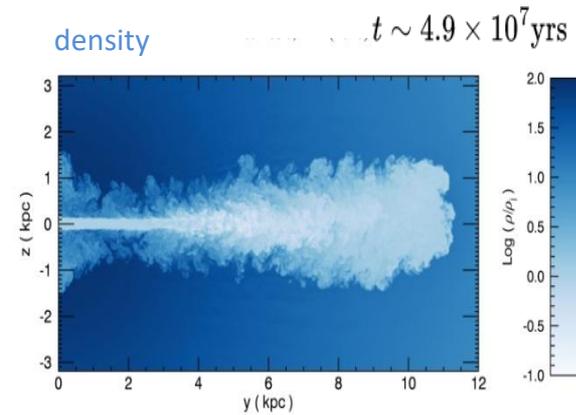
$$\gamma \sim 1 - 2$$

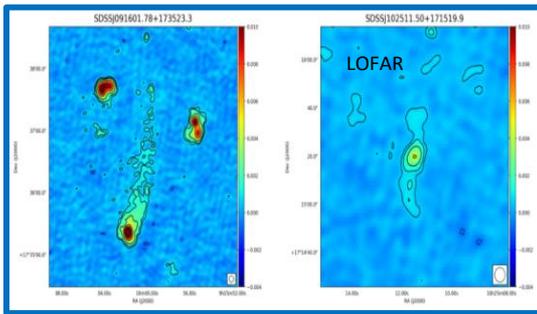
We perform a set of simulations with similar conditions

$$M = 4 \quad (0.06 c)$$

$$\eta = 10^{-2}$$

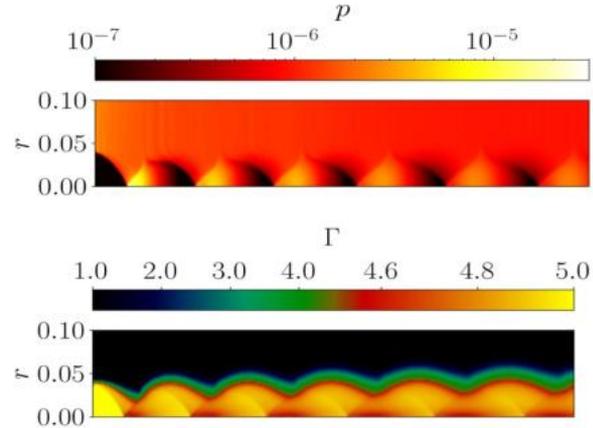
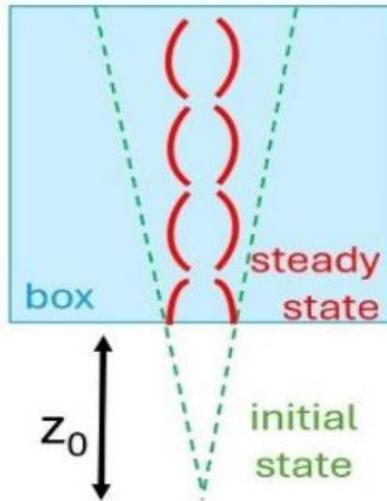
Non-relativistic with 3 different external density stratification (Massaglia et al. 2016)





For the very low power jet (FR0) can we find a mechanism that stop the Jet at even smaller distance?

The external pressure or the wind surrounding the jet may collimate the jet inducing a series of recollimation shocks



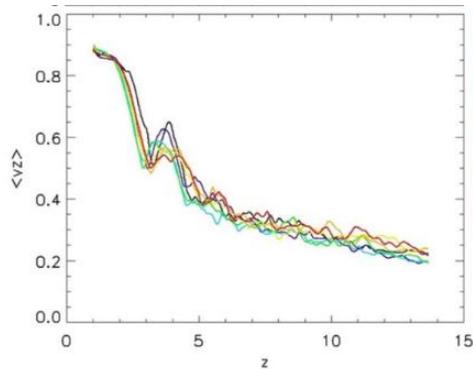
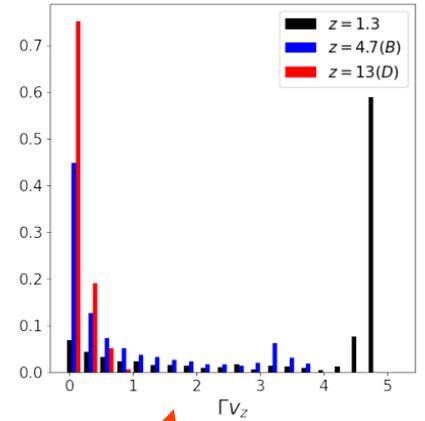
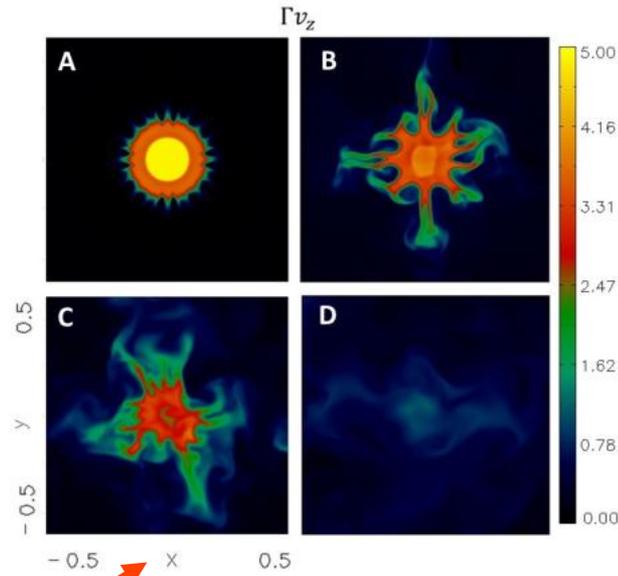
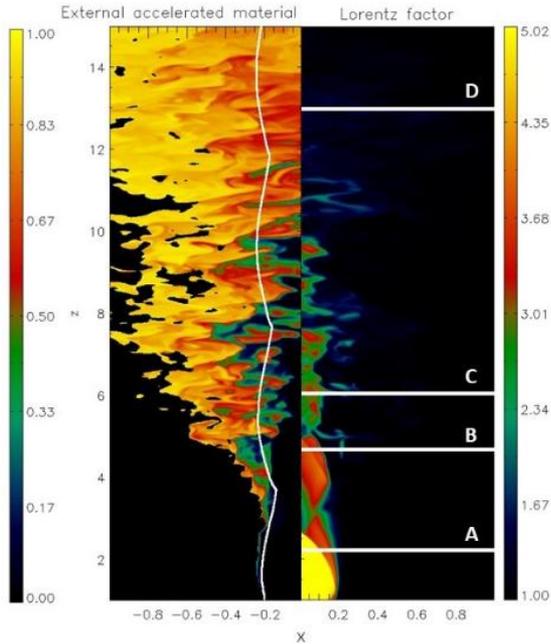
$$\gamma = 5 \quad v = 7.6 \times 10^{-5} \quad L_j = 6.5 \times 10^{39} \text{ erg s}^{-1}$$

$$kT/mc^2 = 3.9 \times 10^{-5}$$

2D studies of recollimation process show a formation of a chain of recollimation shocks, but in 3D this structure results to be unstable. The mechanism is the centrifugal inst. Connected with the curvature of the fluid streamlines.

# RECOLLIMATION SHOCKS AND FRO

Power  $10^{40}$  erg  $s^{-1}$



Entrainment

Relativistic ( $\gamma = 5$ ) up to about 3pc  
 Mildly relativistic up to about 10 pc  
 Subrelativistic at larger distances

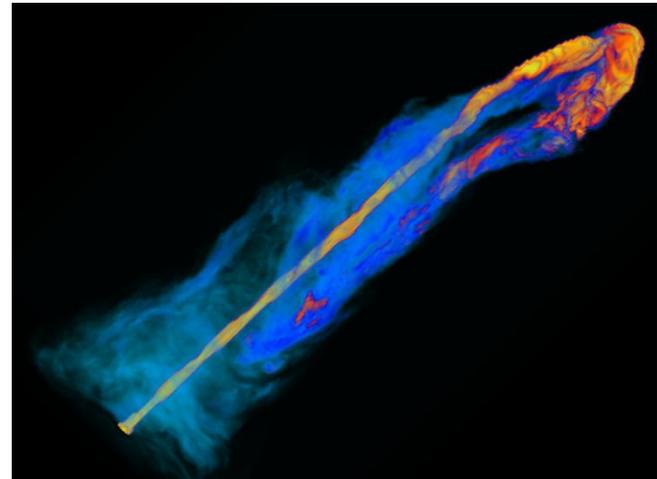
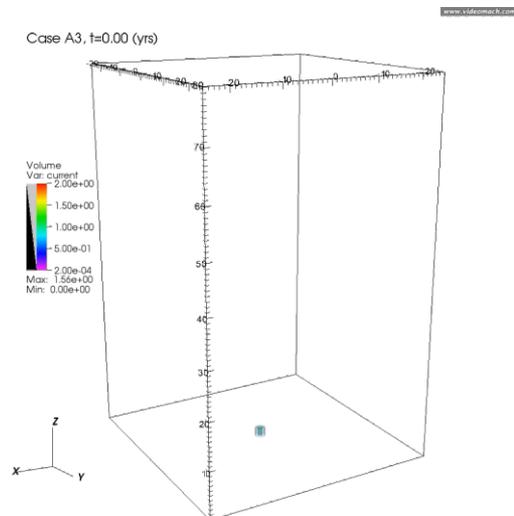
Quasi steady-state

Costa+, 2024,2025

## What about Current-Driven instabilities?

**Dissipation & Particles acceleration:** from the acceleration process we expect that the toroidal magnetic field component is important → this may lead to current driven instabilities → formation of current sheets and magnetic reconnection region → Particles acceleration

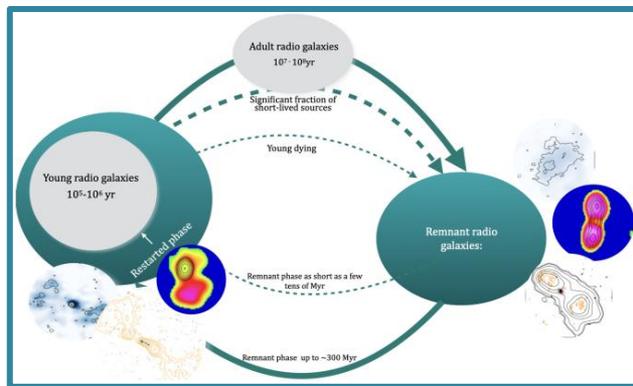
**Jet morphologies:** 3D simulations on a relativistic jet with a toroidal magnetic field show a complex dynamics with bendings and wigglings of the jet head.



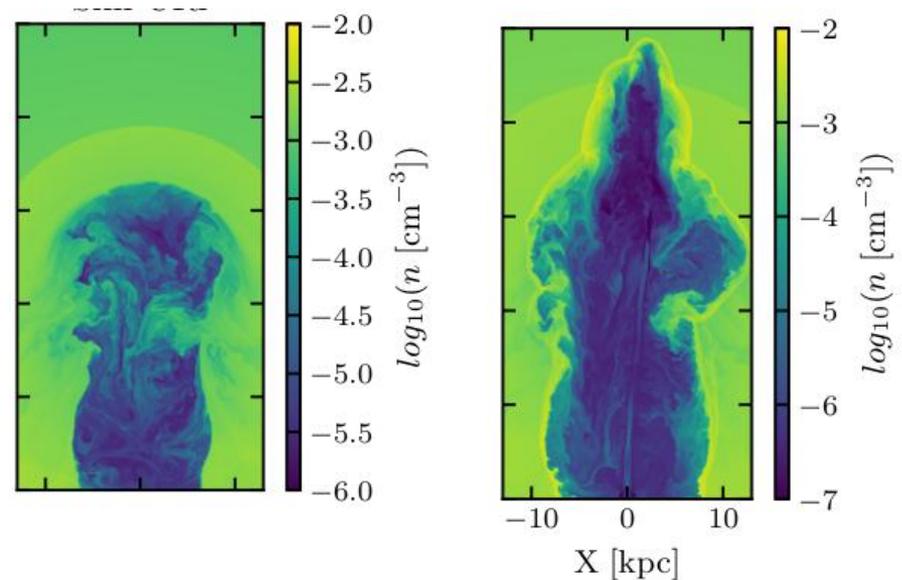
# Termination phase

Jets deposit their energy either at the head (FR II) or gradually along their trajectory (FR I). this process has an impact on the ambient medium and may have consequences both for the host galaxies and for the intergalactic medium (jet feedback).

There are recent observational evidences that jets can have active recurrent phases



Morganti+ 2023



Ratnaparkhi+ 2026

# From macro to micro physics

## Jets with new hybrid scheme: macro-particle\* + fluid

$$\frac{d\mathbf{x}_p}{dt} = \mathbf{v}(\mathbf{x}_p) \quad \text{Macro-particles move at the fluid velocity}$$

\*to each macro-particle we associate a distribution function of non-thermal relativistic particles

$$\mathcal{N}(p, \tau) = \int d\Omega p^2 f_0 \approx 4\pi p^2 f_0$$

$$\frac{d\mathcal{N}}{d\tau} + \frac{\partial}{\partial E} \left[ \left( -\frac{E}{3} \nabla_\mu u^\mu + \dot{E}_l \right) \mathcal{N} \right] = -\mathcal{N} \nabla_\mu u^\mu \quad \text{Evolution equation for the distribution function}$$

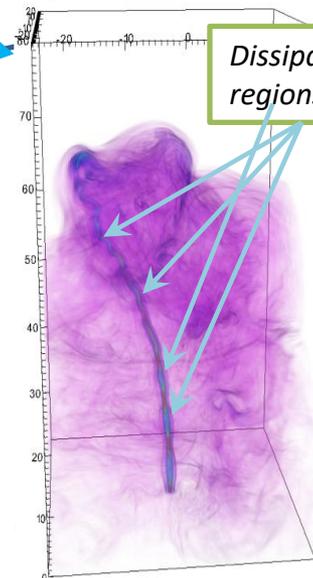
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Energy losses, synchrotron, inverse Compton

Sub-grid model for acceleration at shocks (we take into account shock strength, orientation of magnetic field etc.)

Sub-grid model based on PIC simulations

Transport eq. for cosmic rays Parker 1965, Jokipii & Parker 1970, Skilling 1975, Webb 1989.



Collisionless relativistic shocks

Relativistic magnetic reconnection

Vaydia+ 2018  
Mukherjee+ 2021

# How the RMHD models can talk to radioastronomers

RMHD eqs + particle transport eq



Spatial evolution

They move at fluid velocity

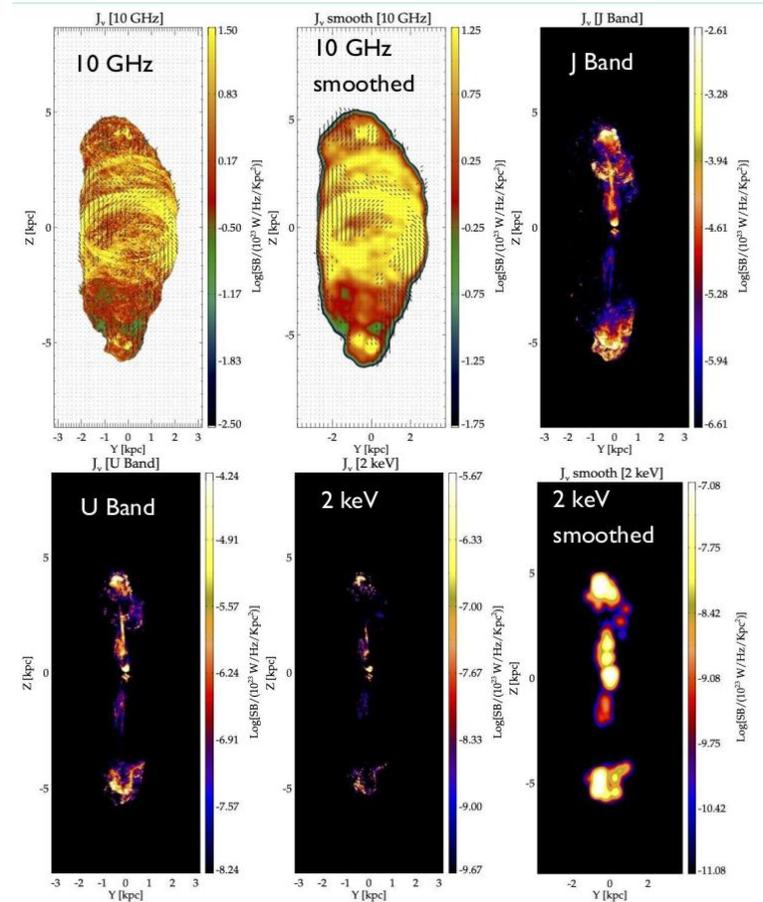
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Spectral evolution

Radiative losses+sub-grid model for acceleration at shocks



Observational properties



## ... next

Bridging the gap between dynamical and radiative properties of jets is fundamental for comparing simulations with observations

This represents a challenge because of the huge scale differences.

Continuing to investigate sub-grid models in order to improve and to keep updated with the latest microphysical results.

Developing a sub-grid model for particles acceleration at magnetic reconnection (Nurisso+ 2023, Suriano+ 2026)

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